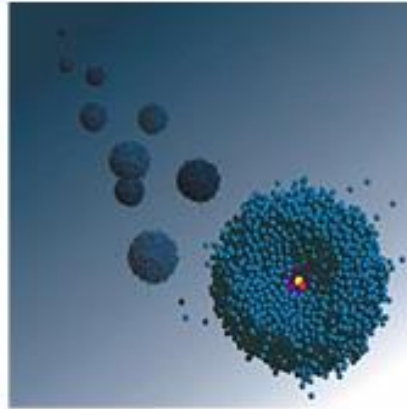


Brioni 2008, 29. August 2008

Are Hydrogen Clusters Superfluid?



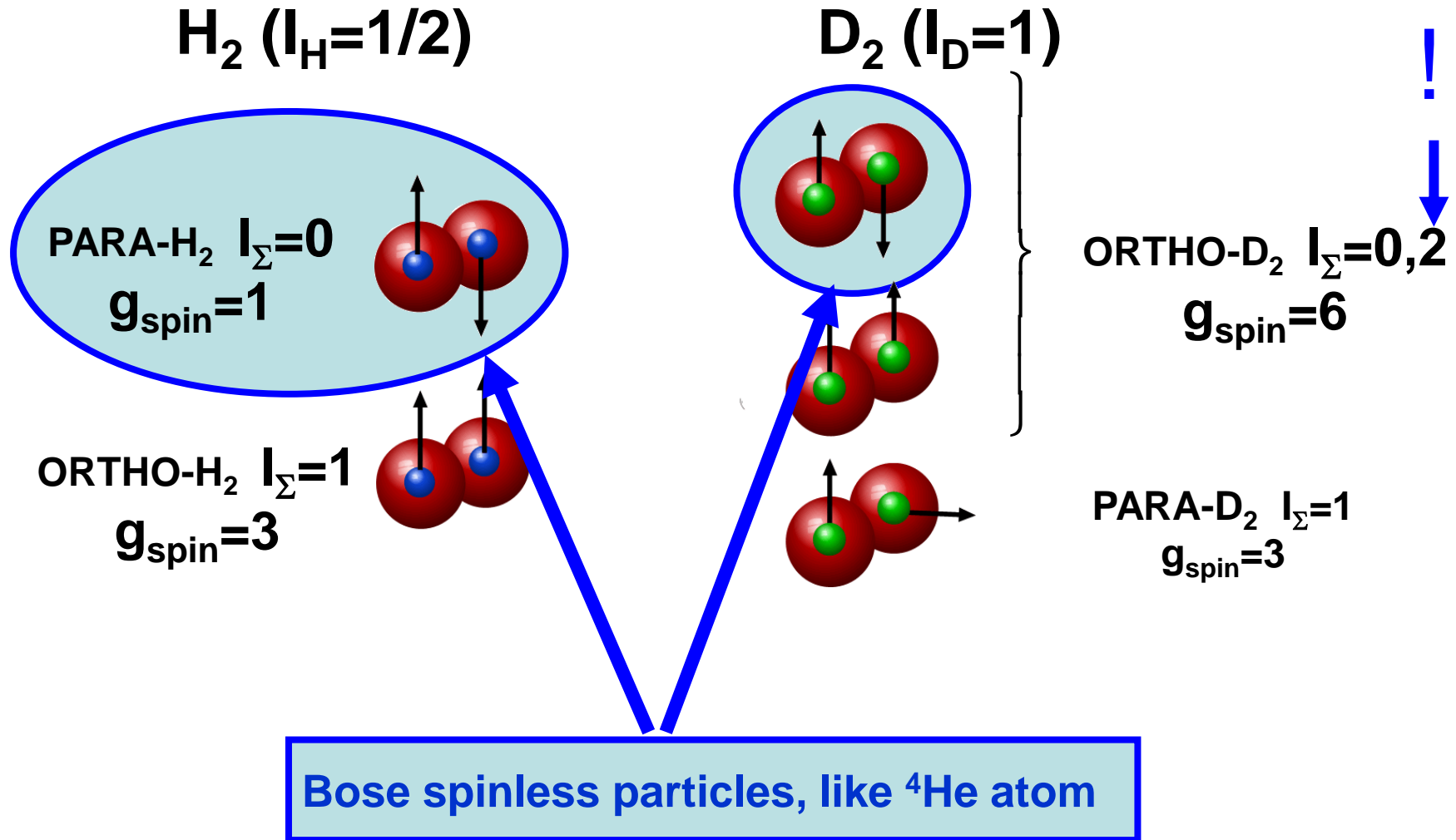
J.Peter Toennies

jtoenni@gwdg.de

Max Planck Institut für Dynamik und Selbstorganisation

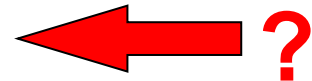
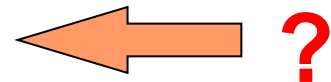
Göttingen, Germany

Hydrogen Isotopomers

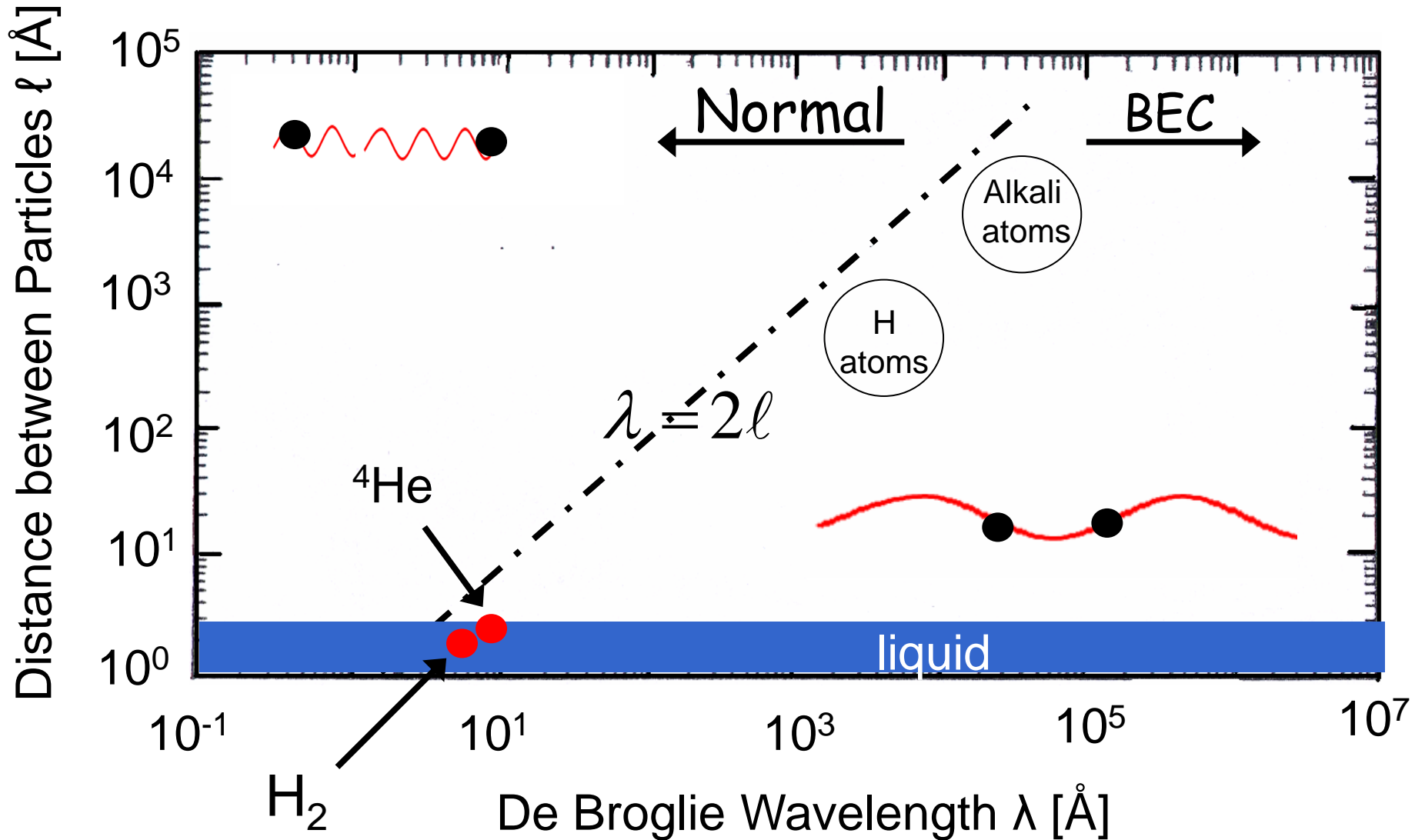


Bosons Have Unique Coherent Macroscopic Quantum States

Particle	Coherence seen in
cooper pairs	superconductivity
alkali gases hydrogen atom gases metastable He gases	BEC
^4He liquid	superfluidity
^4He solid	supersolid
^3He pairs (liquid)	superfluidity
supercooled H_2	superfluidity
nucleon pairing	nuclei
Higgs Boson	elementary particles
Stars	Pulsar glitches



Bose-Einstein Condensation vs. Superfluidity

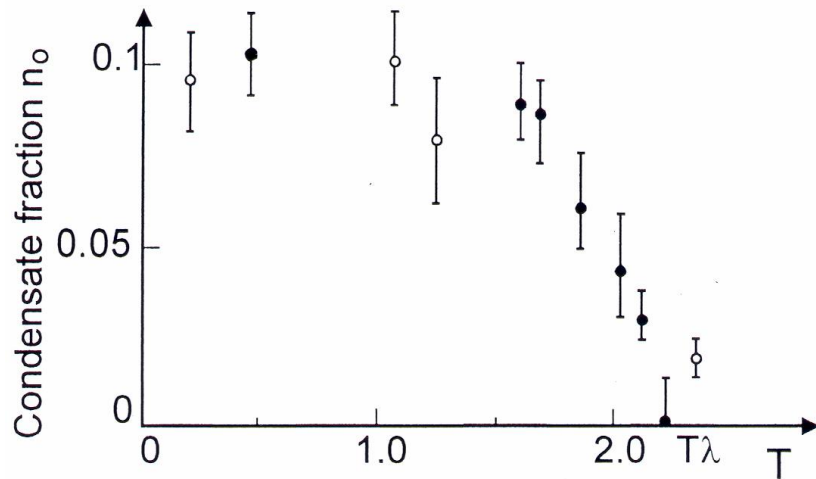


BEC and Superfluidity:

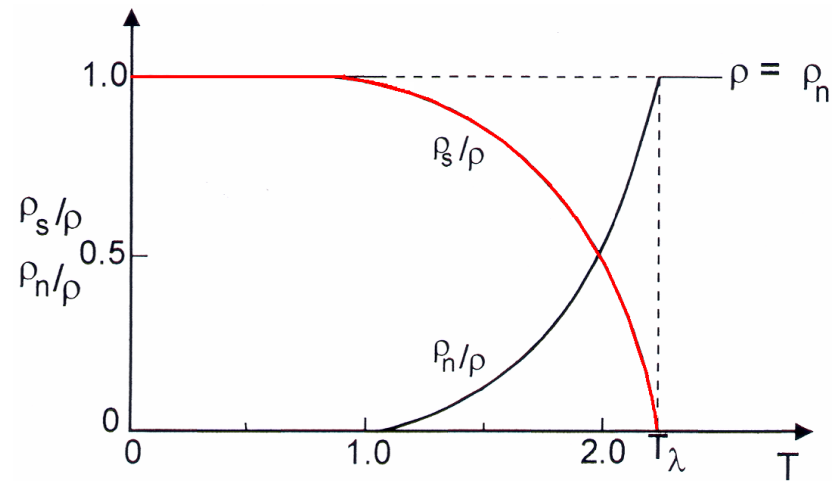
....., the explicit connection between superfluidity and BEC is not trivial and has been the object of a longstanding and deep investigation in the last decades, mainly for its importance in understanding the physics of liquid helium

Dalfovo, Giorgini, Pitaevskii and Stringari: Rev. Mod. Phys. **71** (1991)

In He⁴ Condensate Fraction is only 10%



Superfluid Fraction approaches 100%



Superfluidity is usually regarded as a **macroscopic** phenomenon!!

There are Well Known **Microscopic** Manifestations of Superfluidity

SOVIET PHYSICS JETP

VOLUME 37 (10), NUMBER 1

JANUARY, 1960

SUPERFLUIDITY AND THE MOMENTS OF INERTIA OF NUCLEI

A. B. MIGDAL

Submitted to JETP editor February 13, 1959

J. Exptl. Theoret. Phys. (U.S.S.R.) 37, 249-263 (July, 1959)

A method is developed for the treatment of superfluidity of nuclei. A formula which agrees satisfactorily with experiment is obtained for the moment of inertia of a nucleus. An expression is found for the change in the energy of "pairing" in the transition from an even-even to an even-odd nucleus, and also for the change in the moment of inertia associated with this transition.

The existence of correlated pairs and superfluidity is evidenced most clearly in nuclear moments of inertia. The moments of inertia of nuclei are two or three times smaller than those computed from the formula for the moment of inertia of a solid, and this is the most direct evidence for the superfluidity of nuclear matter.

Search for Superfluidity in para-H₂

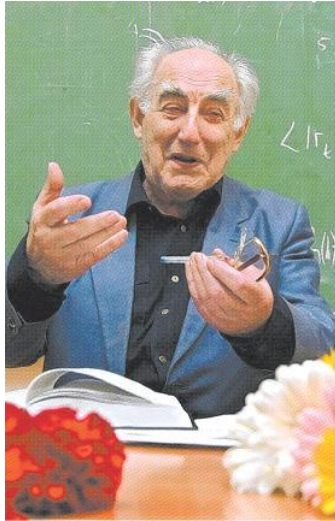
In 1972 Vitali Ginzburg (Nobel Prize 2003) and Alexander Sobyenin predicted that liquid hydrogen can become superfluid.

Ginzburg and Sobyenin, JETP Lett. 15, 242 (1972)

ideal gas BEC transition temperature

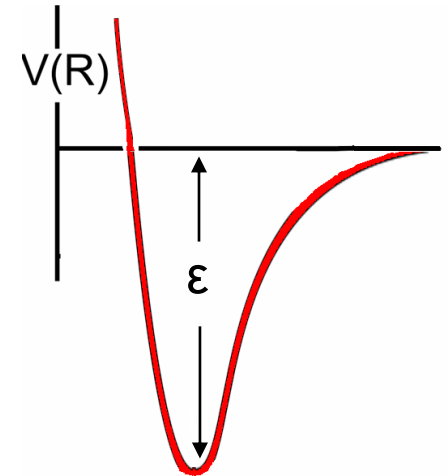
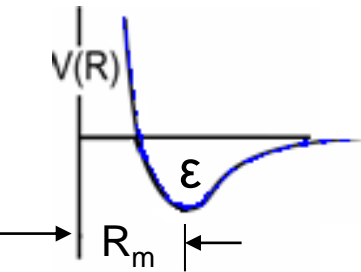
$$T_c = \frac{3.3 \hbar^2}{m k_B} \left(\frac{n}{g} \right)^{2/3}$$

$$T_c(pH_2) = 6.0 \text{ K}$$



Hydrogen is More Strongly Bound and Less Delocalized than Helium

	He	H ₂
R _m [Å]	3.0	3.44
ε [K]	11.0	34.2
Particle vol. [Å ³]	28.9	38.2
Λ*	2.9	2.0



Quantum delocalization is given by the **de Boer parameter**

$$\Lambda^* = \frac{\lambda(\varepsilon)}{R_m} \quad \text{where} \quad \lambda(\varepsilon) = \frac{h}{\sqrt{m\varepsilon}}$$

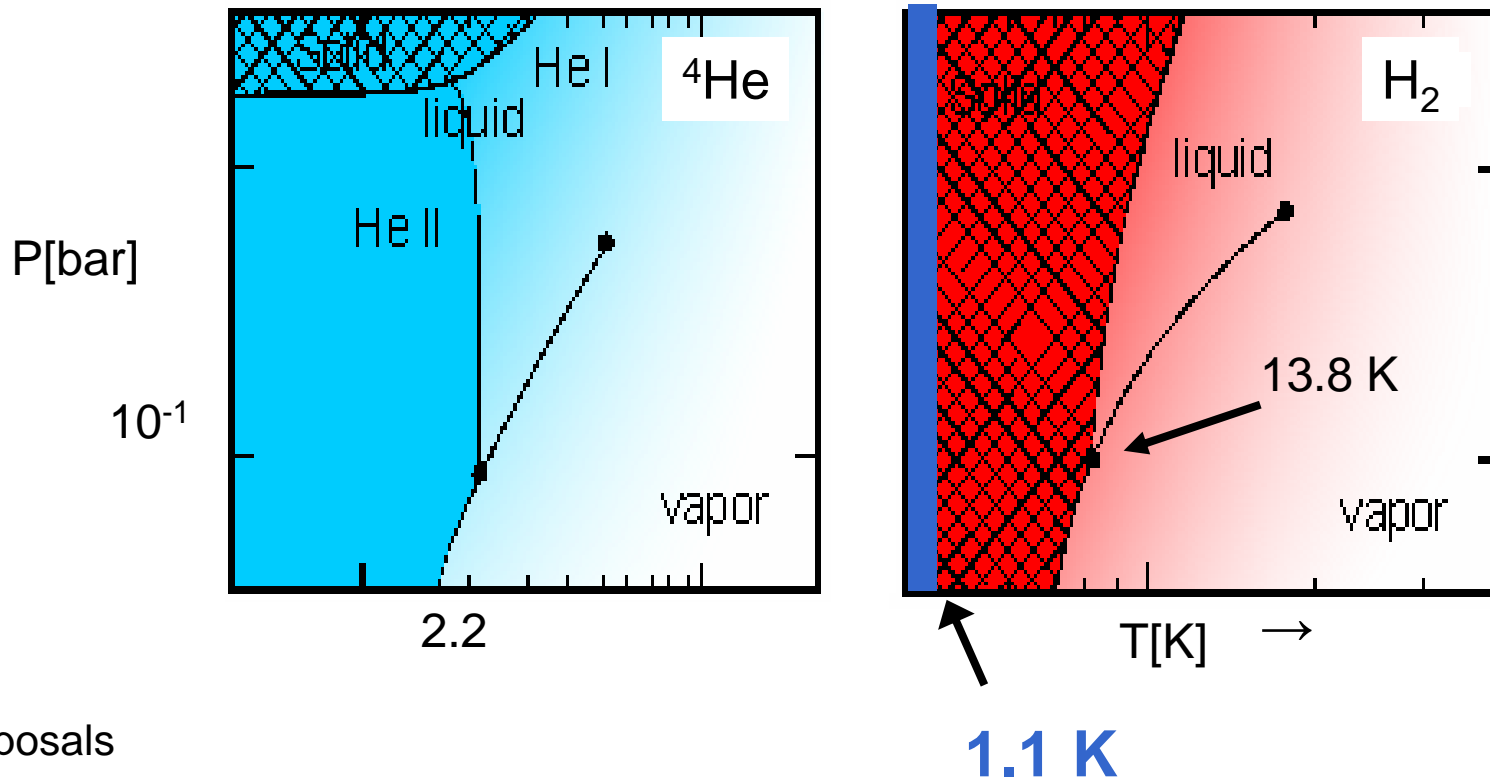
Phase diagram

The Superfluid transition temperature is estimated to be only

$$T_{SF}(pH_2) \approx 1.1K$$

Apenko, Phys. Rev. B60, 3052 (1999)

The Problem with Hydrogen is that it Solidifies at 13.8 K



Methods Proposed for Creating Superfluid H₂

Depression of melting point by going to negative pressures	Theory
Depression of Debye temperature by adding impurities	Theory
Supercooling drops inside liquid ⁴ He	Expt.
Depression of melting point in porous Vycor glass	Expt.
Supercooling droplets in free jet expansions	Expt.
Forced liquid-solid phase equilibrium	Theory
Forced cooling of a microfilament of liquid	Expt.
Small clusters	Theory and Expt.

Methods Proposed for Creating Bulk Superfluid H₂

Depression of melting point by going to negative pressures	Theory	Ginzburg and Sobyenin, 1972
Depression of Debye temperature by adding impurities	Theory	Gelikman, 1974
Supercooling drops inside liquid ⁴ He	Expt.	Maris et al., 1983 – 87
Depression of melting point in porous Vycor glass	Expt.	Torii, Maris and Seidel, 1990
Supercooling droplets in free jet expansions	Expt.	Knuth, Schünemann and Toennies, 1995
Forced liquid-solid phase equilibrium	Theory	Vorob'ev and Malyshenko, 2000
Forced cooling of a microfilament of liquid	Expt.	Grisenti et al., 2006
Small clusters	Theory	Ceperley et al, 1991

Outline

- I. Introduction: BEC and Superfluidity
- II. Theory of Small para-H₂ (pH₂) Clusters
- III. Raman Spectroscopy of Pure pH₂ clusters
- IV. Matter-Wave Diffraction of H₂ Clusters
- V. Spectroscopic Evidence for Superfluidity

Experimentalists

Andrej Vilesov , now USC
Slava Grebenev, now Kirov
Anton Kalinin, Göttingen
Oleg Kornilov, now Berkeley

Theoreticians

David Ceperley, Illinois
Boris Sartakov, Moscow
Mikhail Sevryuk, Moscow

Path Integral Monte Carlo Method I

Partition Function $Z(\beta) = \text{Trace} \exp(-\beta H) = \int dR \rho(R, R; \beta)$

where $\beta = 1/kT$ and $\rho(R, R; \beta)$ is the position space density matrix.

It is a function of the $3N$ coordinates of all the particles.

Via a convolution theorem it can be decomposed into M segments

which is represented by an integral over a path with time steps $\tau = \beta / M$

$$\rho(R_0, R_M; \beta) = \int \dots \int dR_1 dR_2 \dots dR_{M-1} \rho(R_0, R_1; \tau) \rho(R_1, R_2; \tau) \dots \rho(R_{M-1}, R_M; \tau)$$

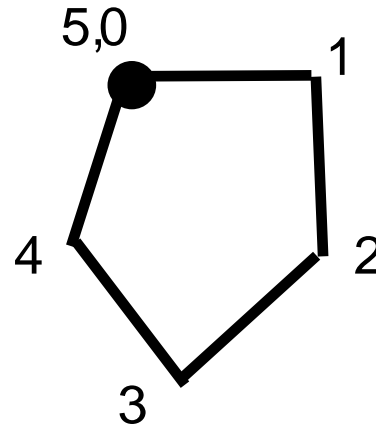
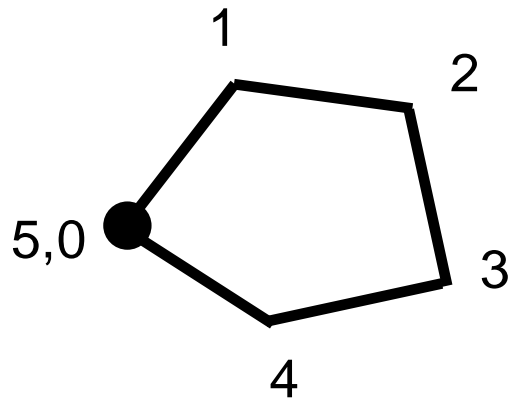
The path for each particle is called a *polymer* and consists of „beads“ and the beads in the different polymers are coupled by the intermolecular potential.

Path Integral Monte Carlo Method II

The size L of each polymer is $L \approx \lambda$

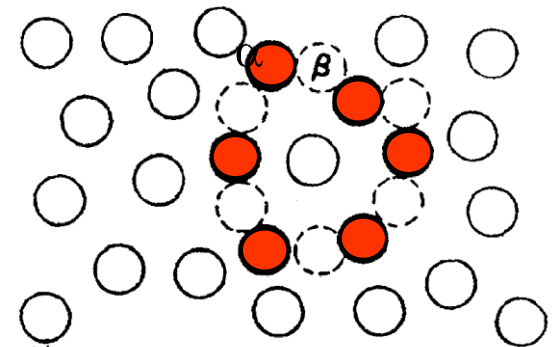
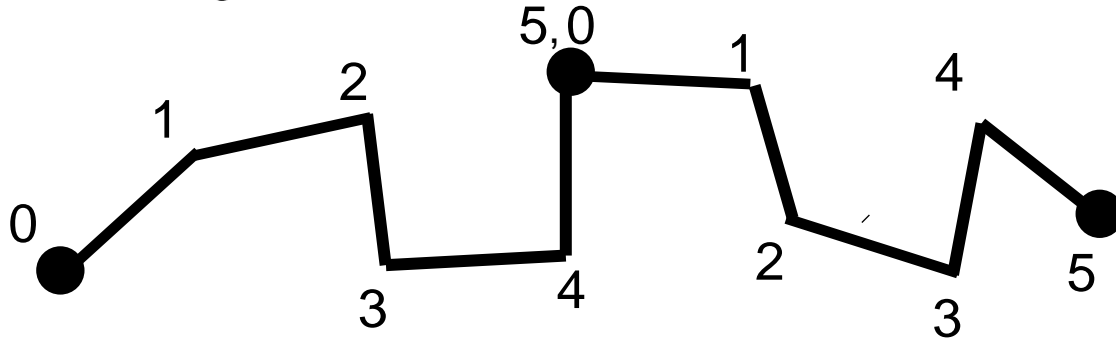
$T > T_C$ $\lambda < l$

separate „polymers“



$T < T_C$ $\lambda \approx l$

linked „polymers“ allow for bose exchange

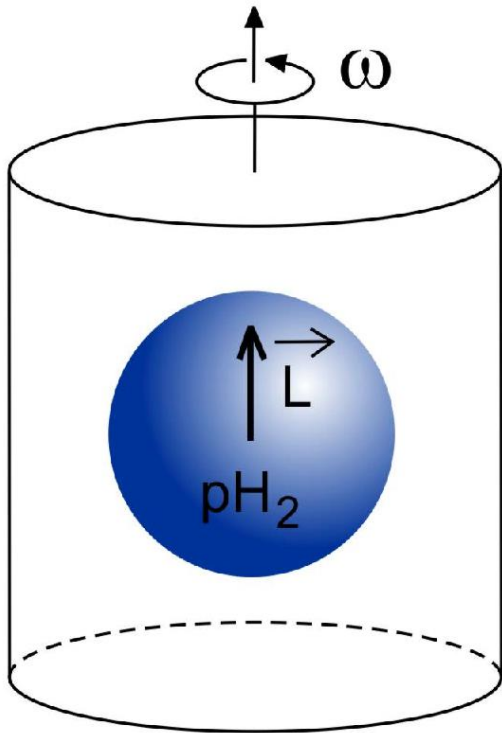


The trajectories are then used to calculate the partition function

Evaluation of Superfluid Response to Rotations: Non-Classical Rotational Moment of Inertia I

Slow rotation of external field with frequency $\vec{\omega}$

$$I_{ij} = \left. \frac{\partial \langle L_i \rangle}{\partial \omega_j} \right|_{\omega=0} = I_{ij}^{cl} - \frac{4m^2 \langle A_i A_j \rangle}{\beta \hbar^2}$$

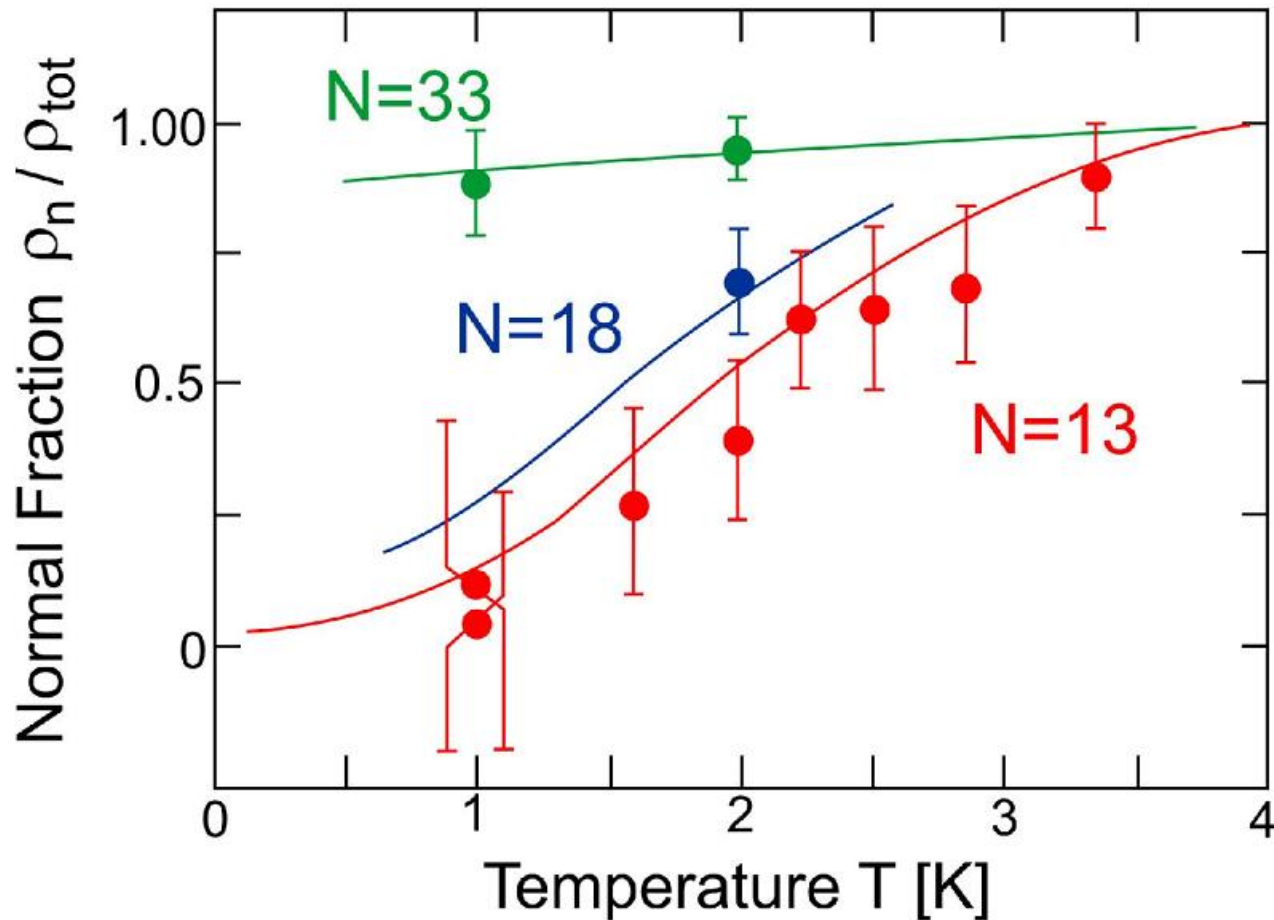


A_i = area of Feynman path projected onto axis i

$$I_{ij}^{cl} = m \int d\vec{r} \rho(\vec{r}) (r^2 \delta_{ij} - x_i x_j)$$

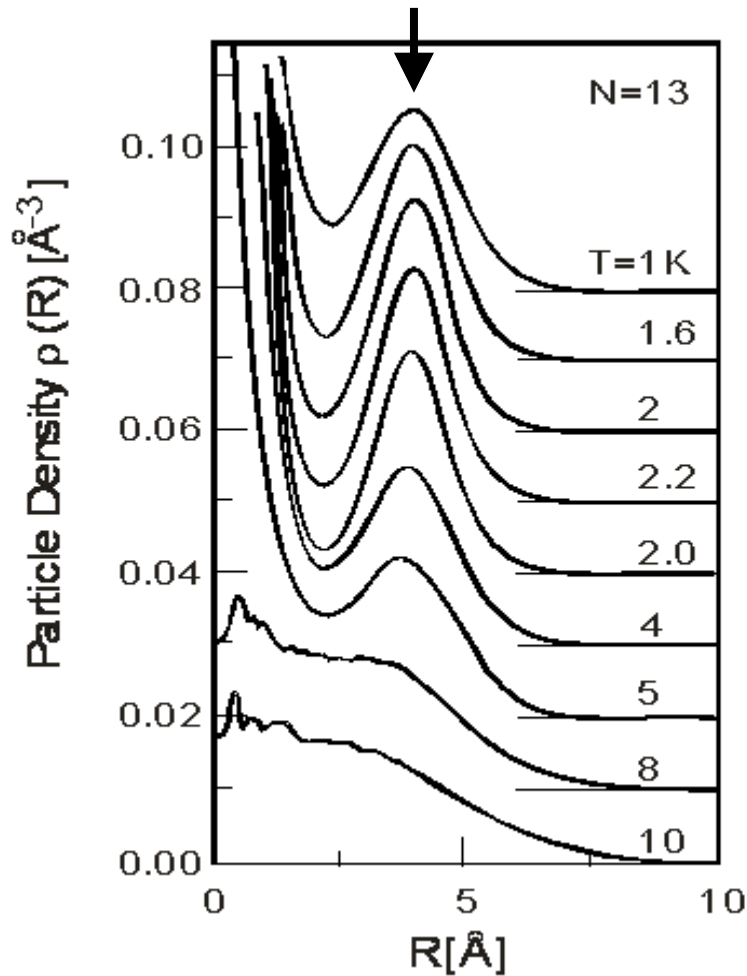
$$f_{\alpha}^s = 1 - \frac{I_{\alpha}}{I_{\alpha}^{cl}} = \frac{4m^2 \langle A_{\alpha}^2 \rangle}{\beta \hbar^2 I_{\alpha}^{cl}}$$

The Reduced Coordination In Small Droplets Favors Superfluid Response



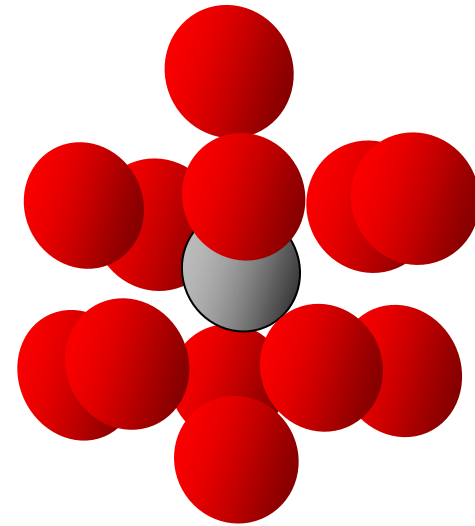
Sindzingre, Klein, and Ceperley, Phys. Rev. Lett. 67, 1871 (1991)

The PIMC Calculated Radial Density Distributions Suggest a „Rigid“-Like Structure



Thermal Melting \rightarrow

Classical $(p\text{-H}_2)_{13}$ Cluster



$N=13$ Icosahedral Structure

Sindzingre, Klein, and Ceperley,

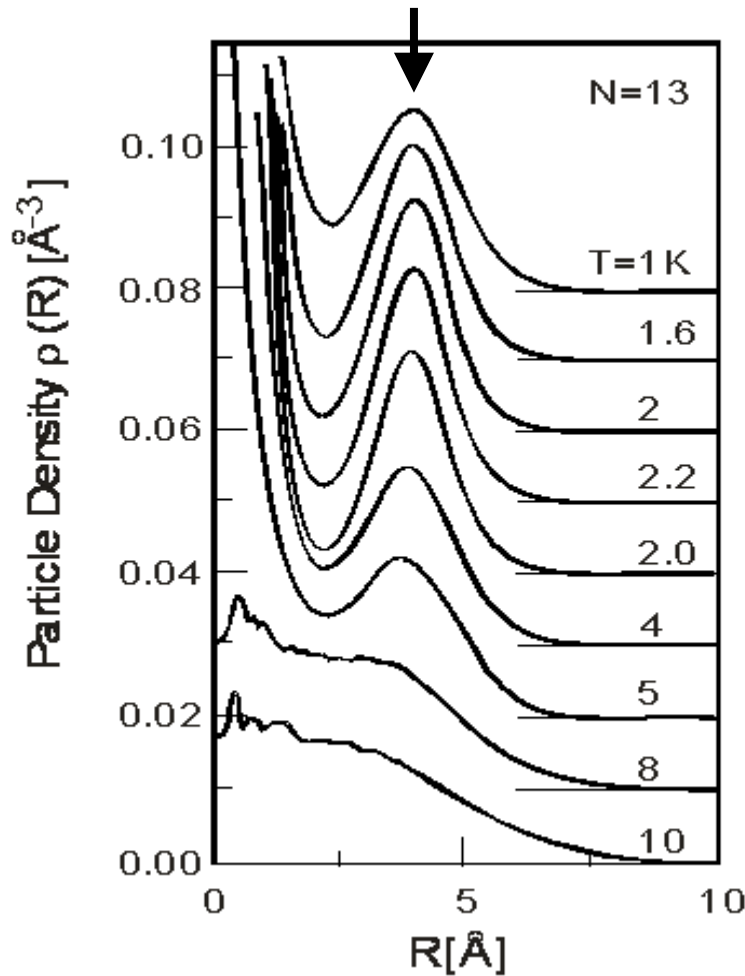
Structure

H_2 : Phys. Rev. Lett. 67, 1871 (1991) ^4He : Phys. Rev. Lett. 63 1601 (1989)

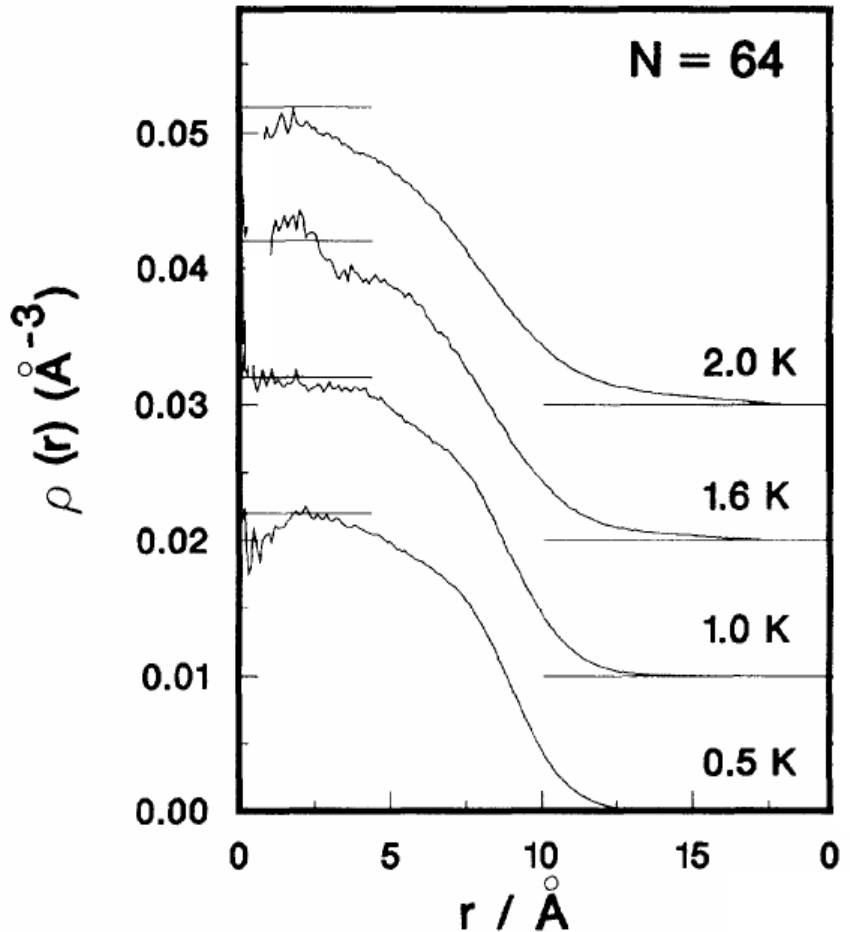
The PIMC Calculated Radial Density Distributions Suggest a „Rigid“-Like Structure

p-H₂ Cluster

Superfluid ⁴He Cluster



Thermal Melting \rightarrow



Sindzingre, Klein, and Ceperley,

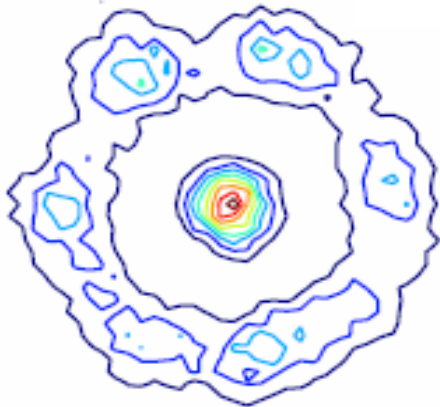
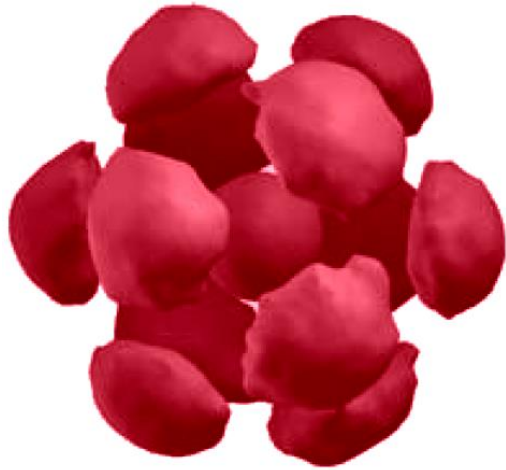
Structure

H₂: Phys. Rev. Lett. 67, 1871 (1991)

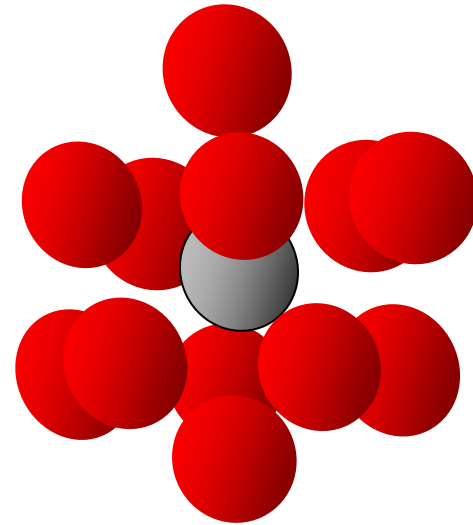
⁴He: Phys. Rev. Lett. 63 1601 (1989)

Structure of a $(p\text{-H}_2)_{13}$ Cluster

Quantum Mechanical

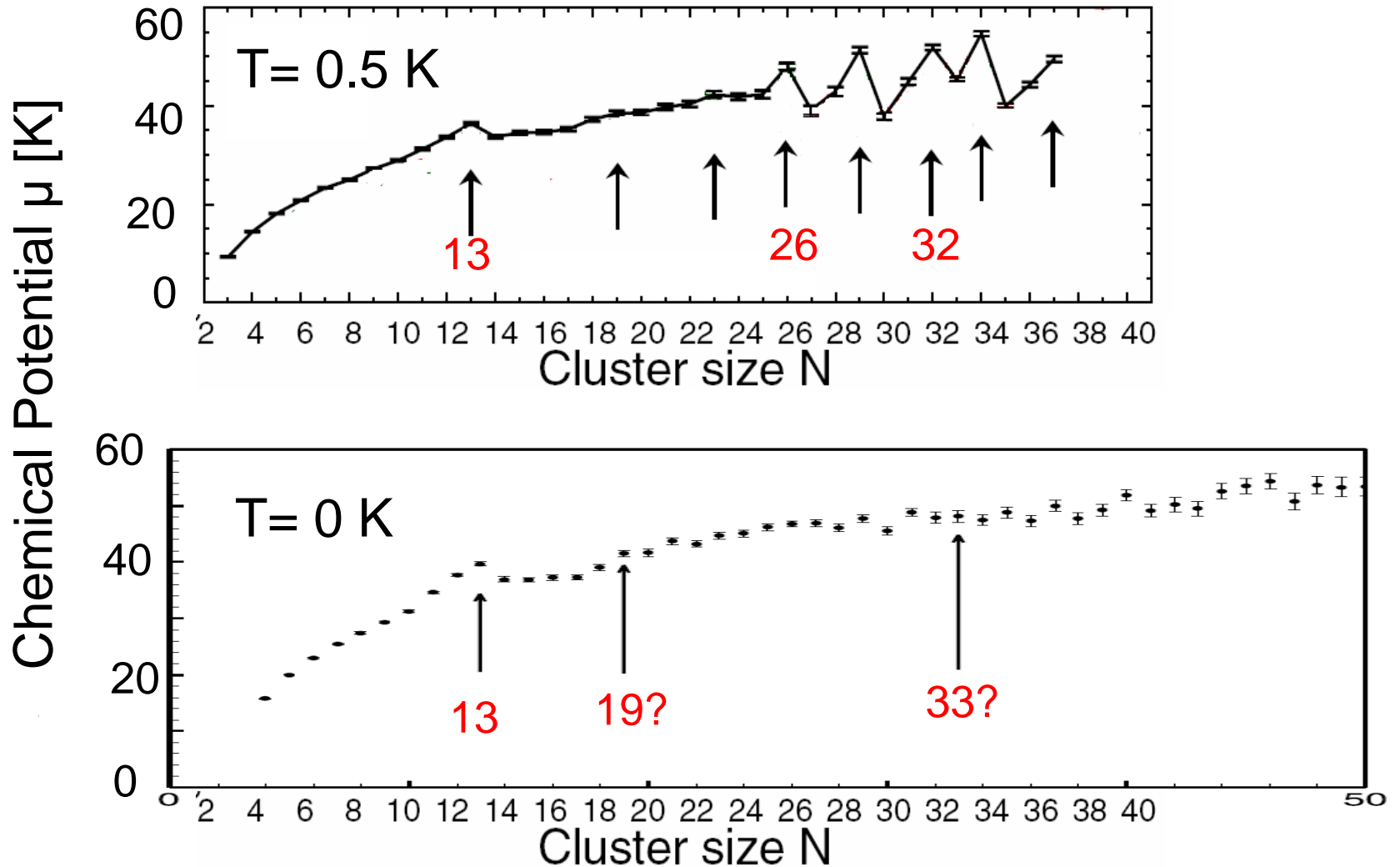


Classical



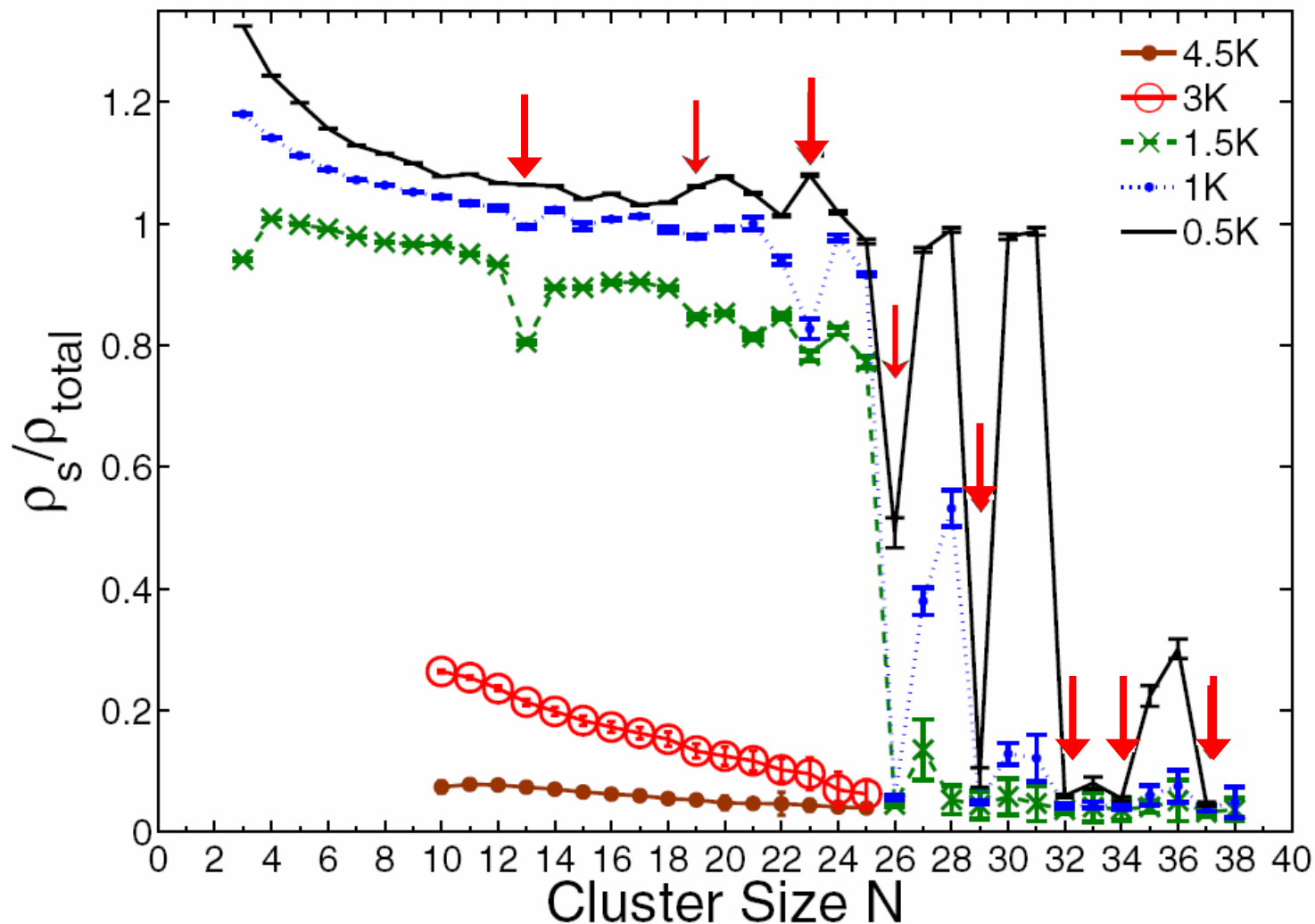
N=13 Icosahedral Structure

The Chemical Potential ($\mu=\Delta E/\Delta N$) Reveals Magic Stabilities



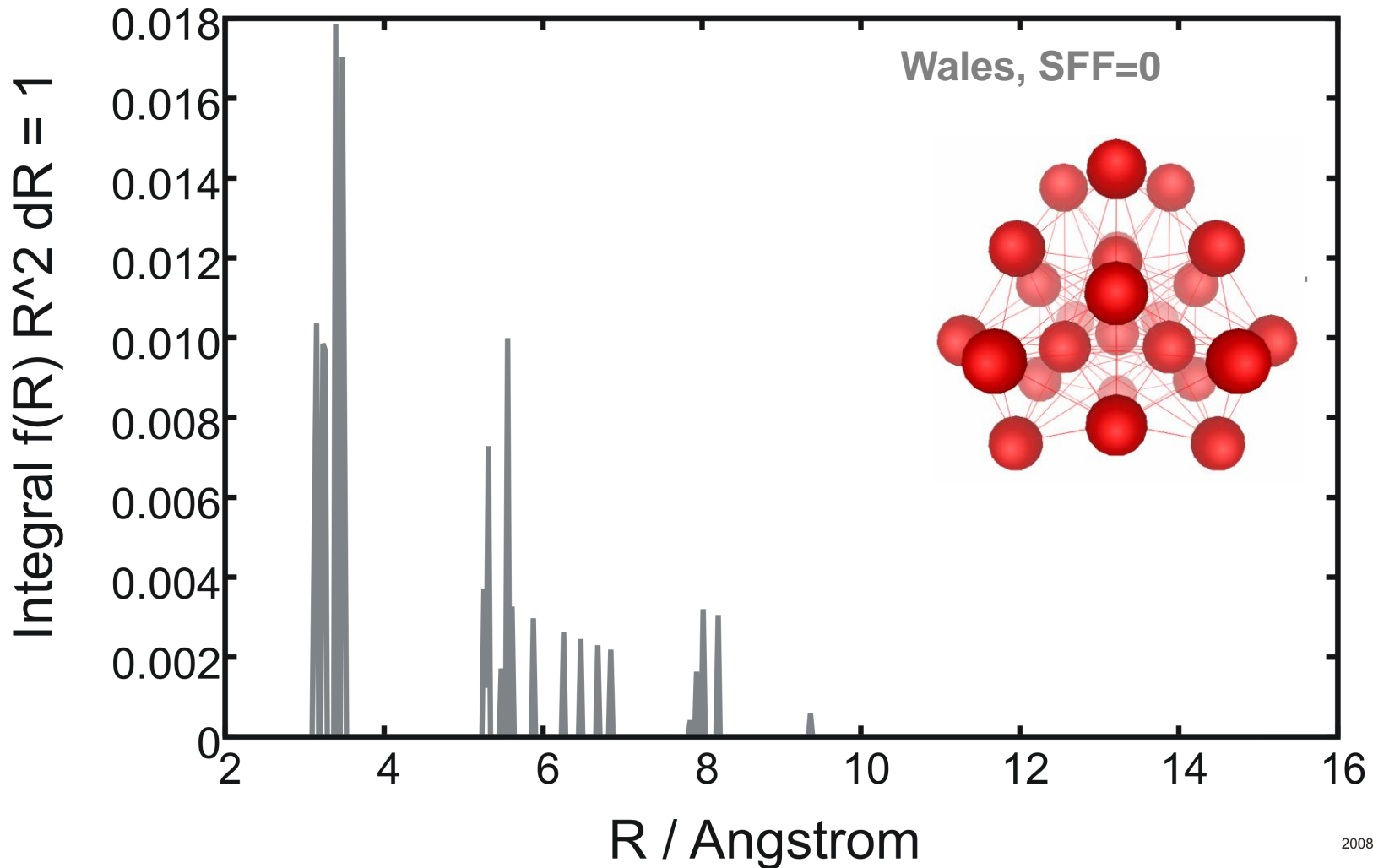
Is the Disappearance at $T=0$ due to Quantum Melting??

The Superfluid Fraction is Depressed at the Magic Numbers



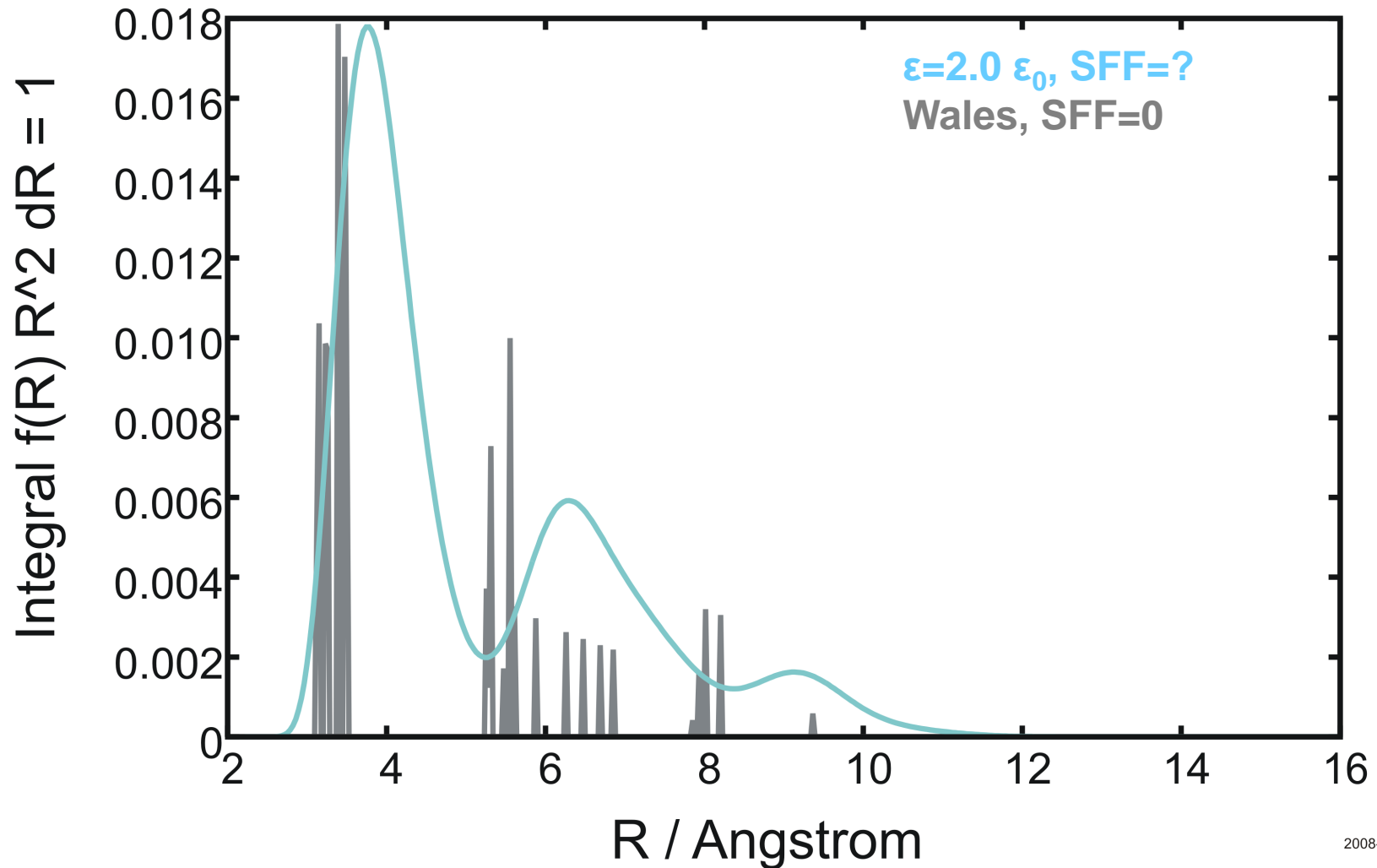
Khairallah, Sevryuk, Ceperley and Toenniesl, Phys. Rev. Lett. 98, 183401 (2007).

By Varying the Potential Well Depth Can Study Transition from Solid to Liquid Cluster (N=26, T=0.5K). Call this „Potential“ Melting



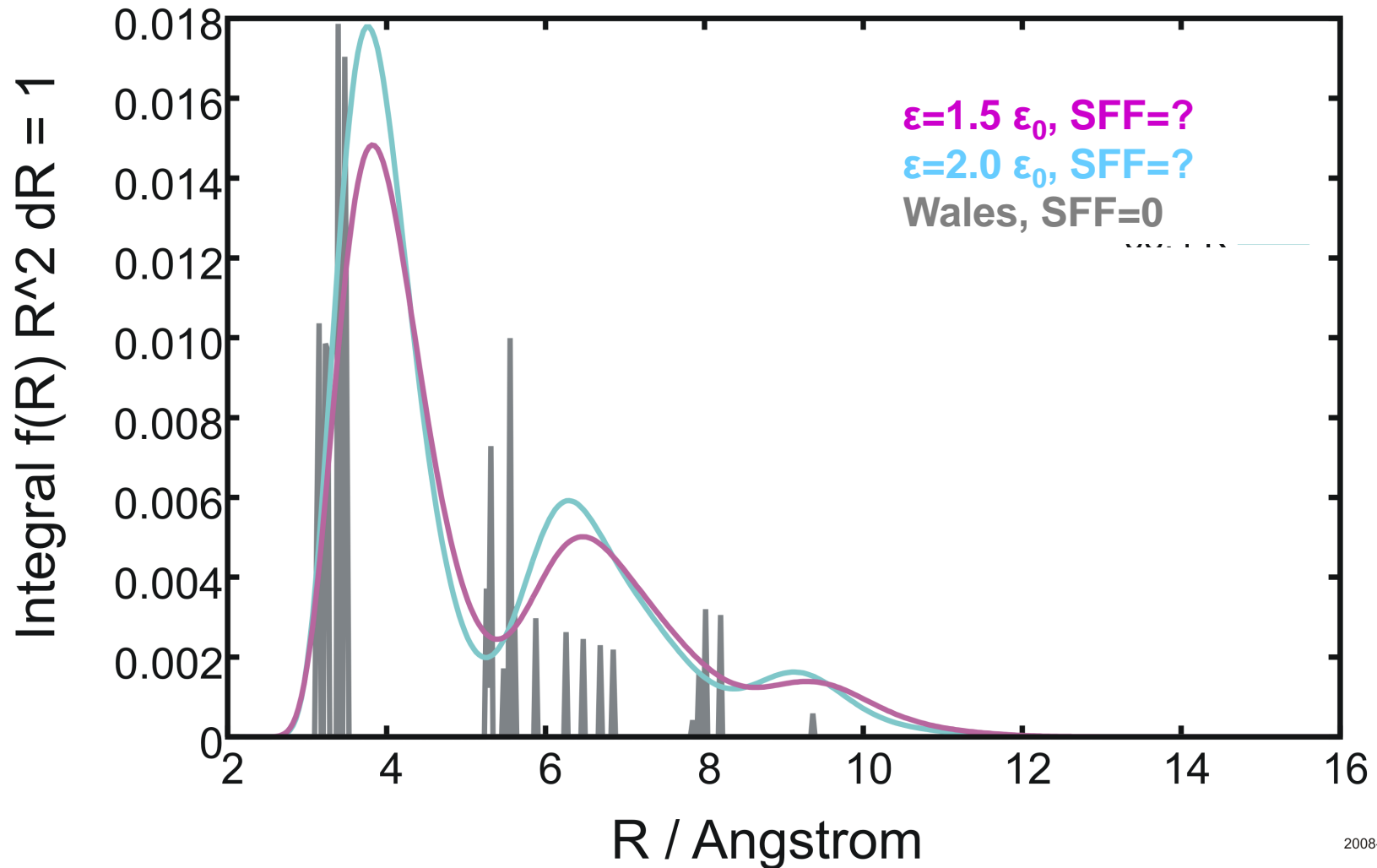
Radial density profiles f for para-clusters: $N = 26$ at $T = 0.5$ K

Various epsilons (34.2 K is the "true" one); the Wales structure scale is arbitrary



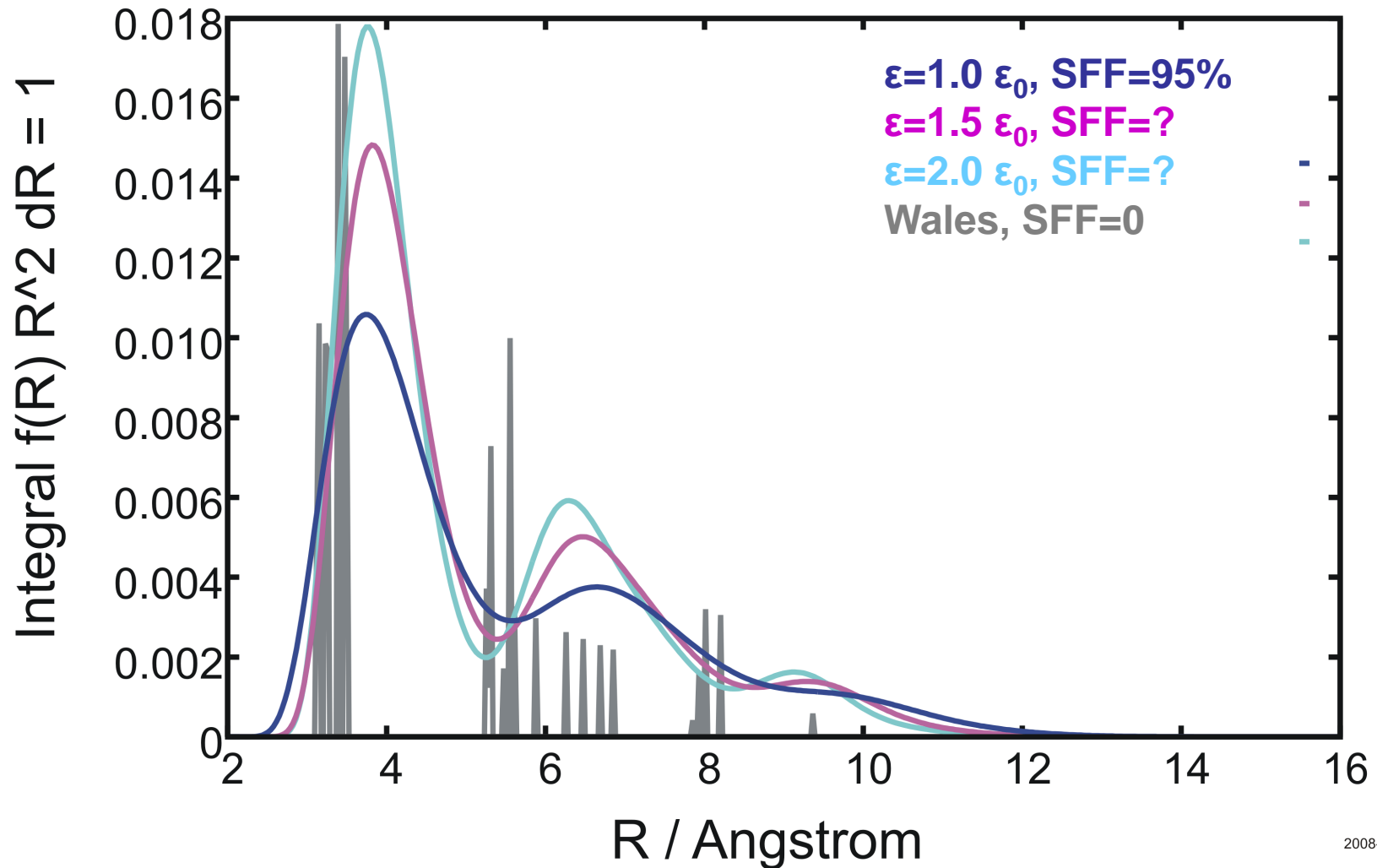
Radial density profiles f for para-clusters: $N = 26$ at $T = 0.5$ K

Various epsilons (34.2 K is the "true" one); the Wales structure scale is arbitrary



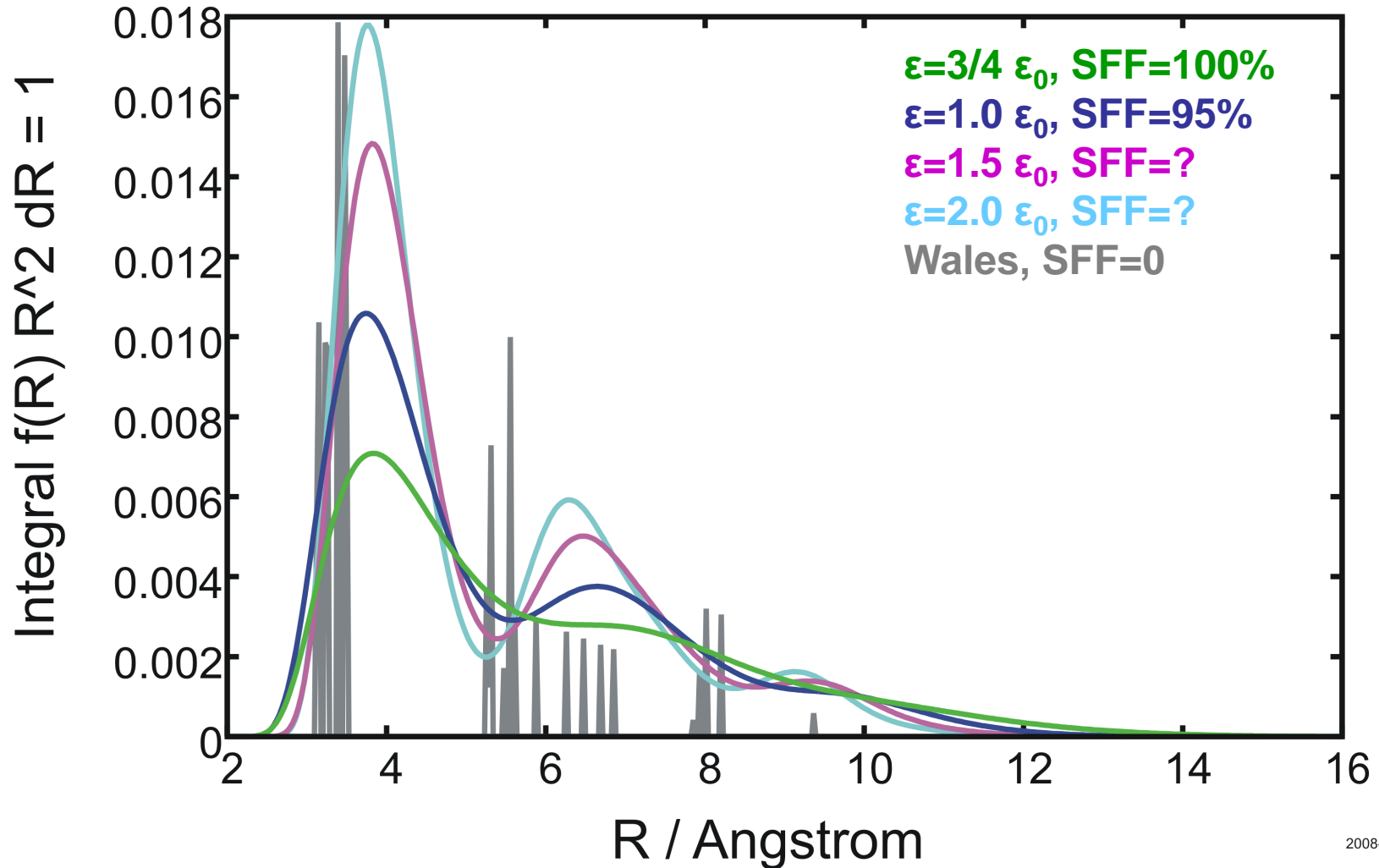
Radial density profiles f for para-clusters: $N = 26$ at $T = 0.5$ K

Various epsilons (34.2 K is the "true" one); the Wales structure scale is arbitrary



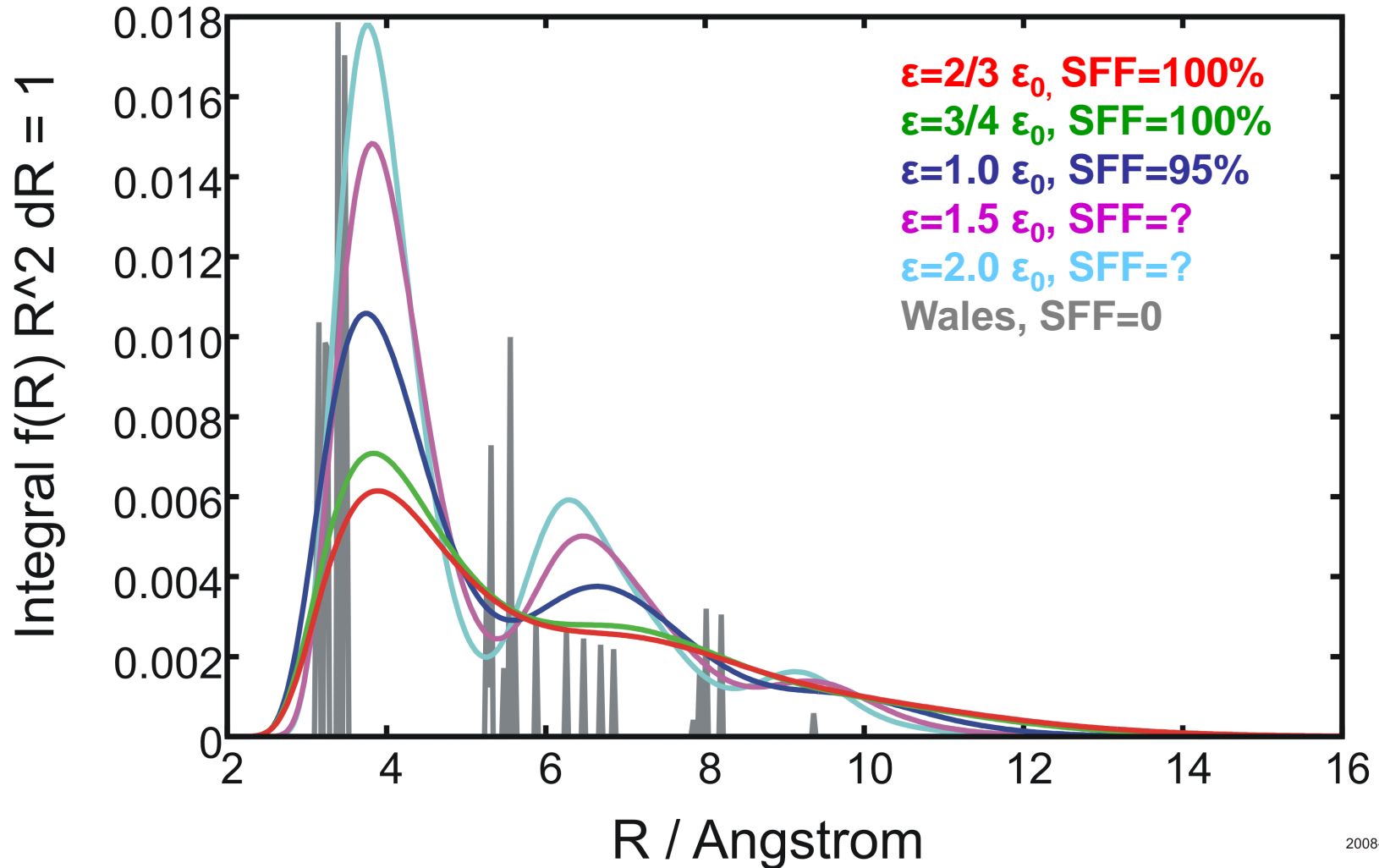
Radial density profiles f for para-clusters: $N = 26$ at $T = 0.5$ K

Various epsilons (34.2 K is the "true" one); the Wales structure scale is arbitrary



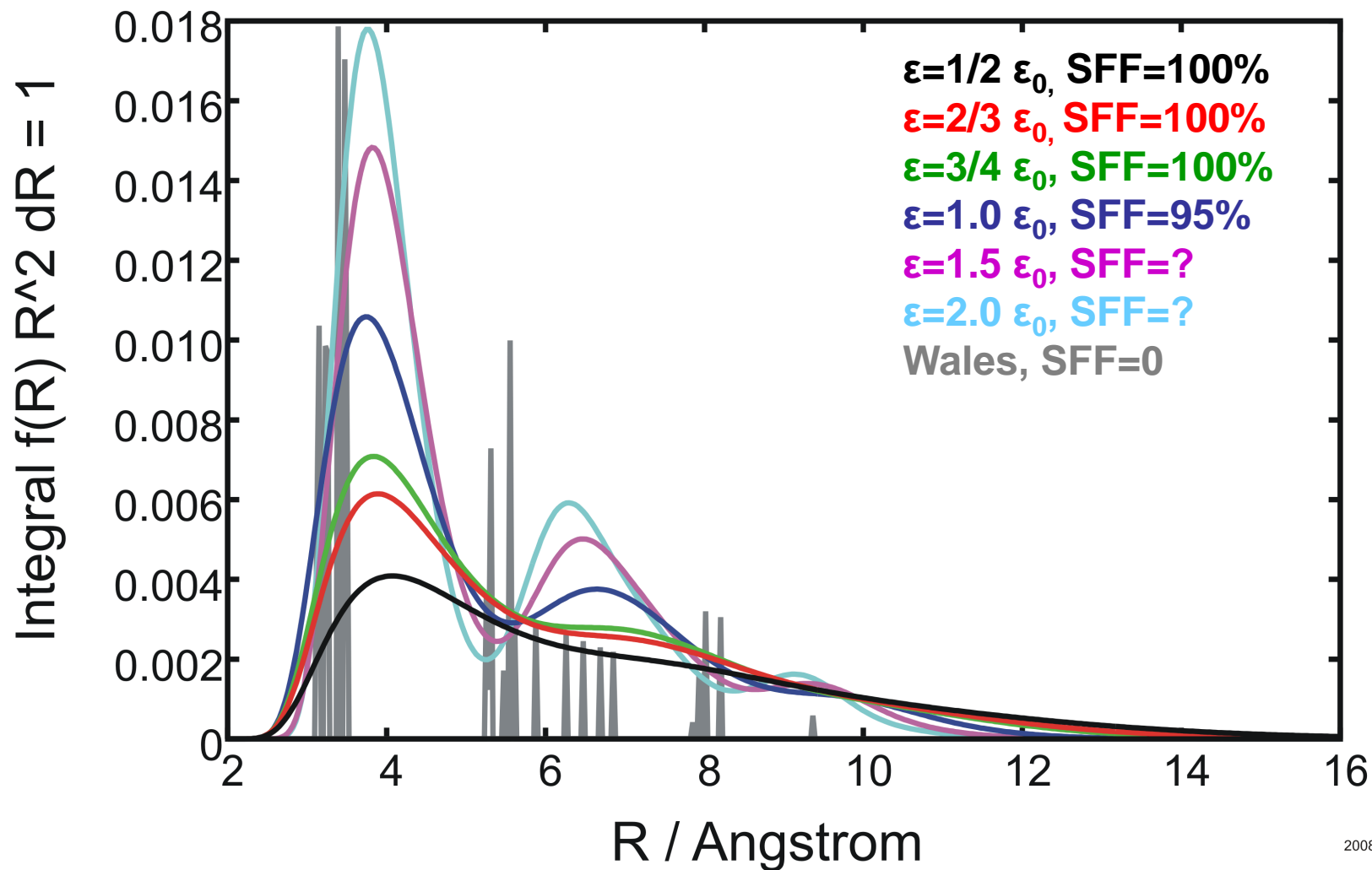
Radial density profiles f for para-clusters: $N = 26$ at $T = 0.5$ K

Various epsilons (34.2 K is the "true" one); the Wales structure scale is arbitrary

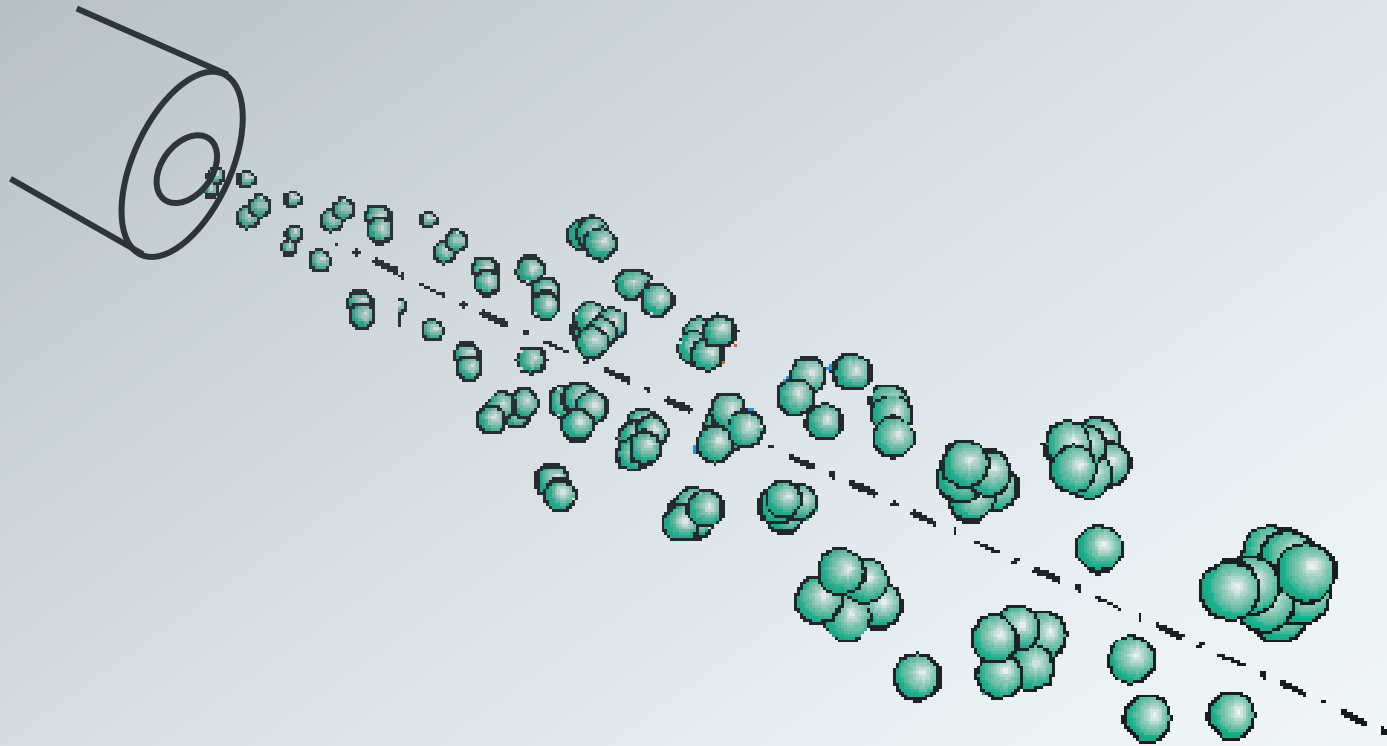


Radial density profiles f for para-clusters: $N = 26$ at $T = 0.5$ K

Various epsilons (34.2 K is the "true" one); the Wales structure scale is arbitrary



Clusters Experiments



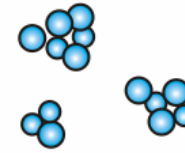
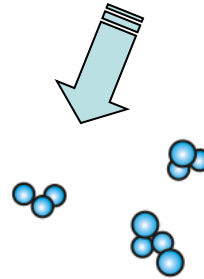
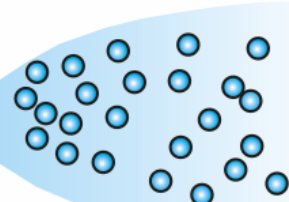
Free Jet Expansions Provide Most Rapid Cooling and Low Evaporation Temperatures

Hetero-nucleation minimized:

purity of gas $< 10^{-4} \rightarrow$

much less than 1 impurity per cluster

$T_0 = 20 \text{ K}, P_0 = 1 \text{ bar}$



**By seeding in He
ambient temperatures
can be reduced to
0.1 - 0.5 K**

Very rapid adiabatic cooling:

$$dT/dt \approx 5 \cdot 10^9 \text{ K sec}^{-1}$$

Melting point is reduced in small droplets

$$T_m/T_{tp} = 1 - \exp(-0.11 \ln N)$$

$$\approx 0.64 \quad (N = 10^4)$$

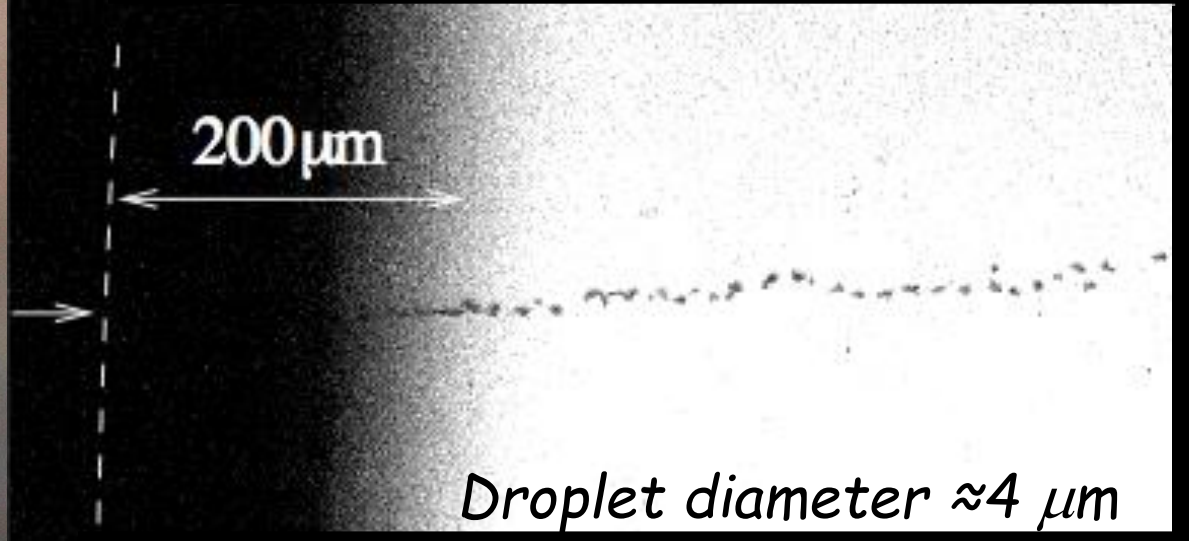
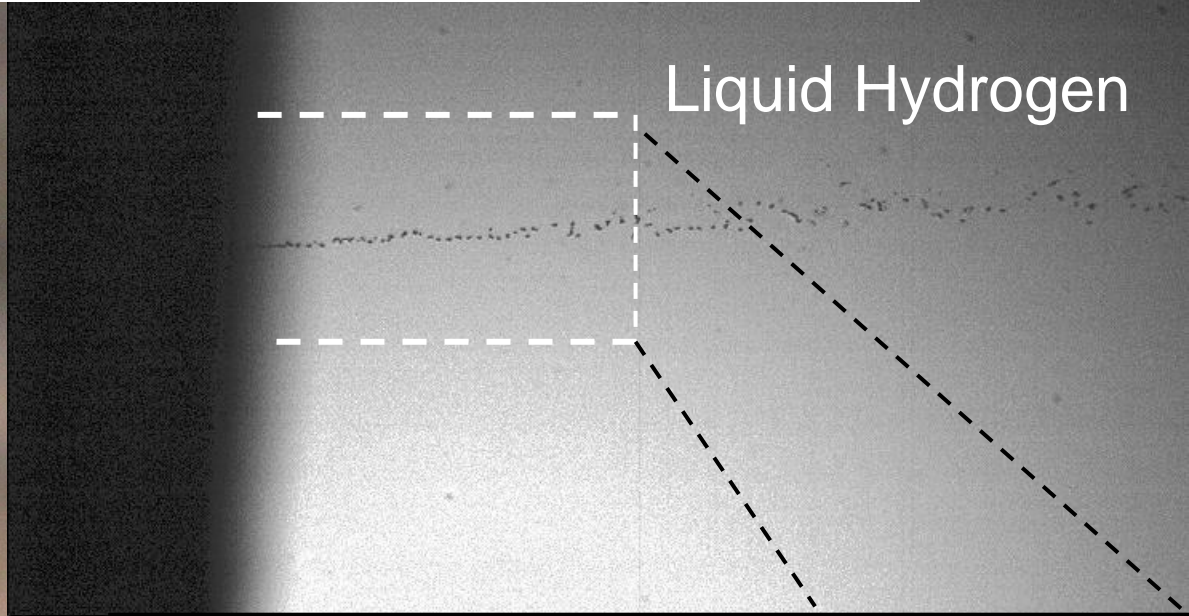
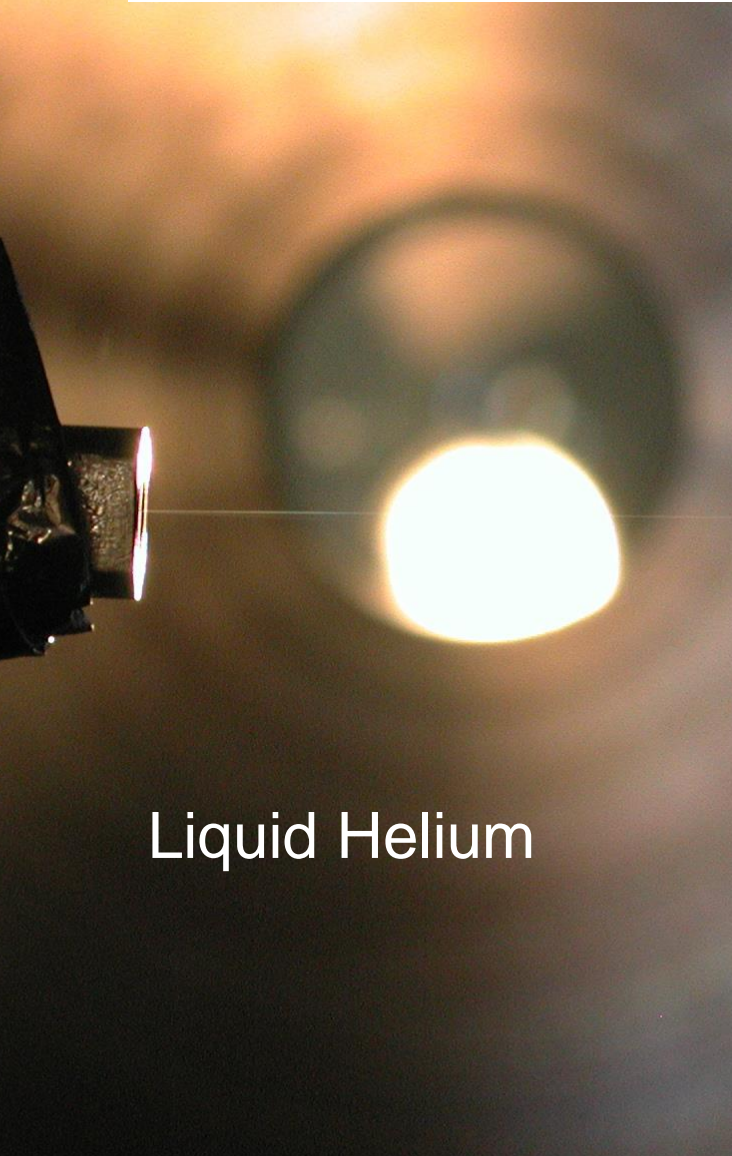
Evaporation cooling

Farges Rule: $T_\infty \propto \Delta h_{evap}$

$$\Delta h_{evap} (H_2) = 109 \text{ K},$$

$$T_\infty = 5.75 \text{ K} \quad (4.5 \text{ K})$$

Rayleigh Breakup of a Liquid Jet



Detection of H₂ Clusters is Problematic

- Spectroscopies

 - ~~Neutron scattering~~

 - ~~Electron scattering~~

 - ~~X-ray scattering~~

 - ~~UV absorption~~

 - ~~Infra-red~~

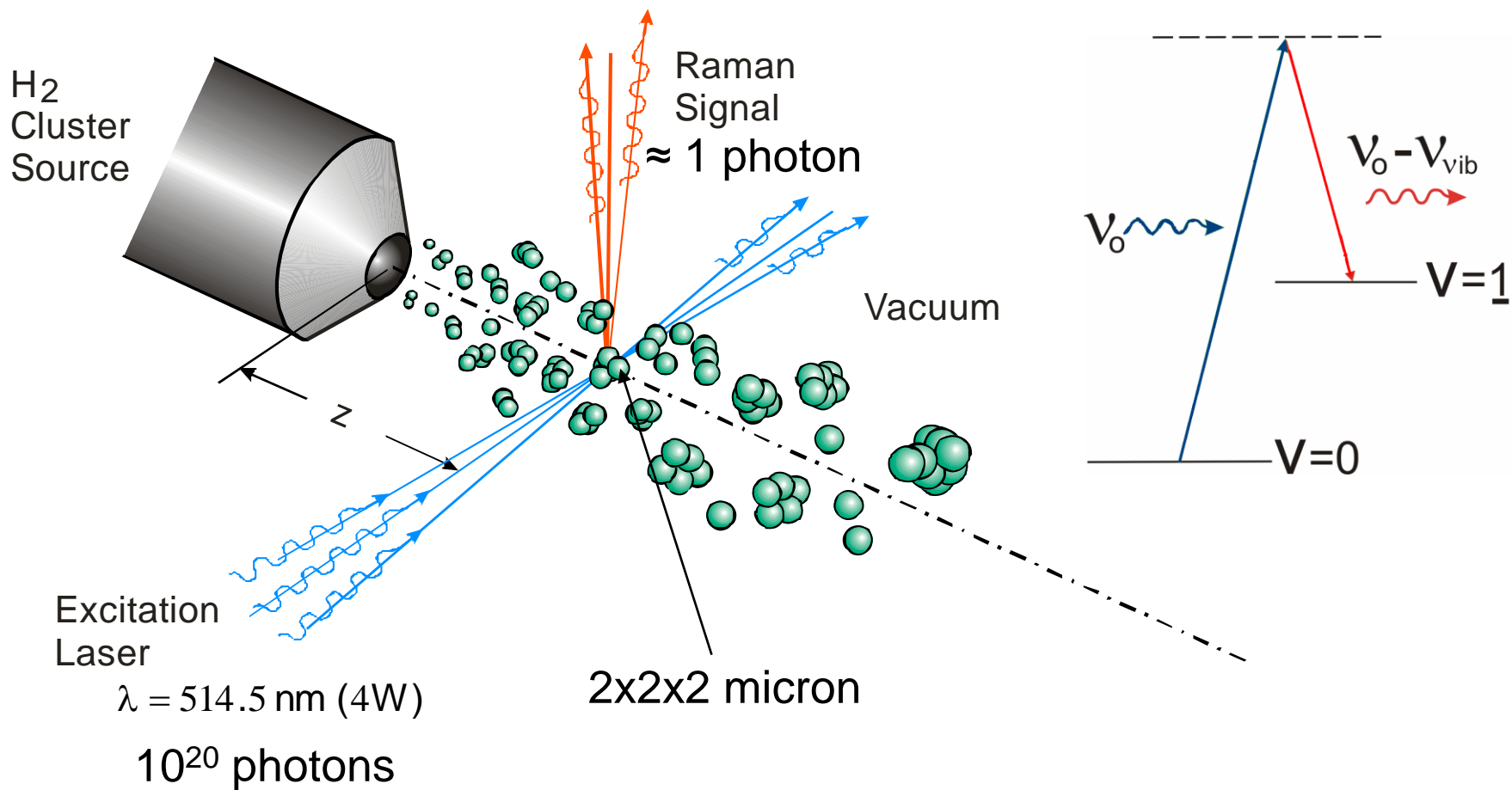
 - Raman scattering

- Others

 - ~~Mass spectroscopy~~

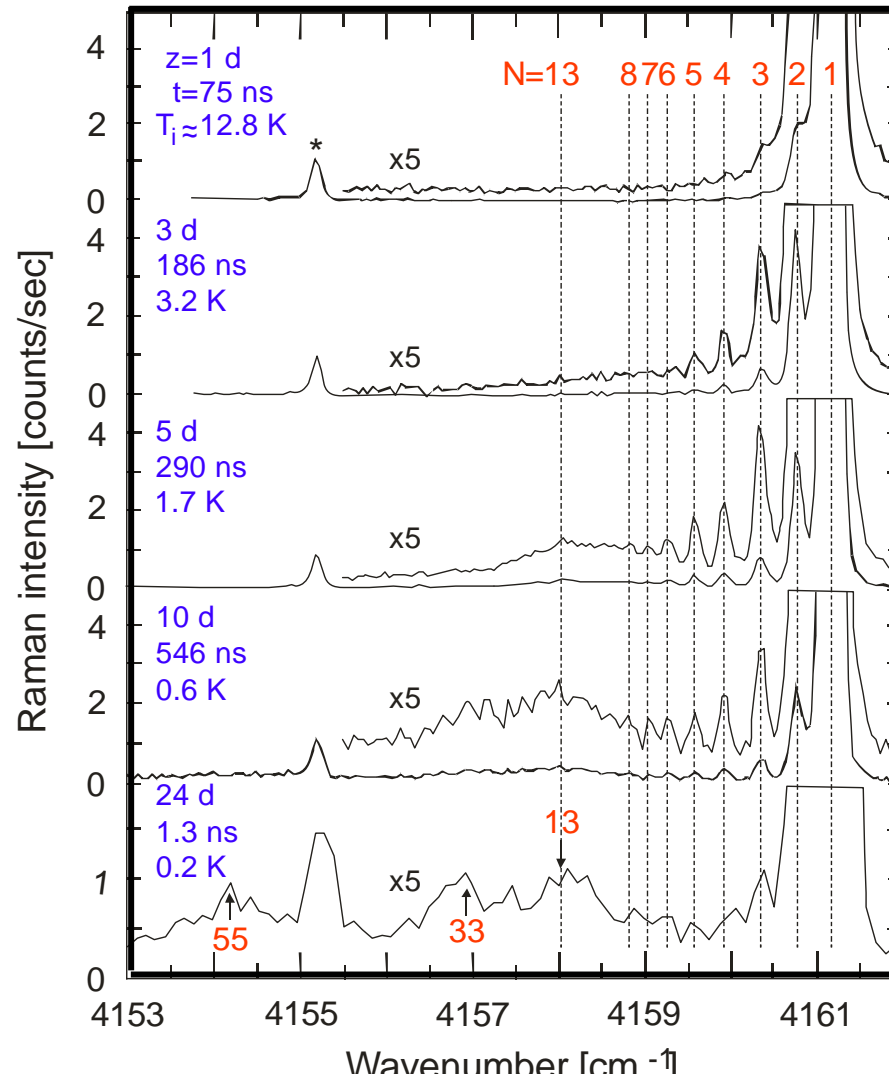
 - Matter-wave diffraction

Raman „Microscopic“ Spectroscopy Inside Free Jet Expansions: Small Clusters of p-H₂



2004-04-19-T4-Ba

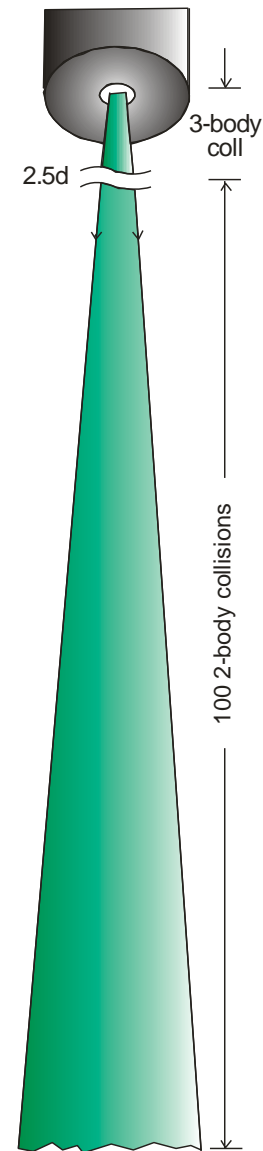
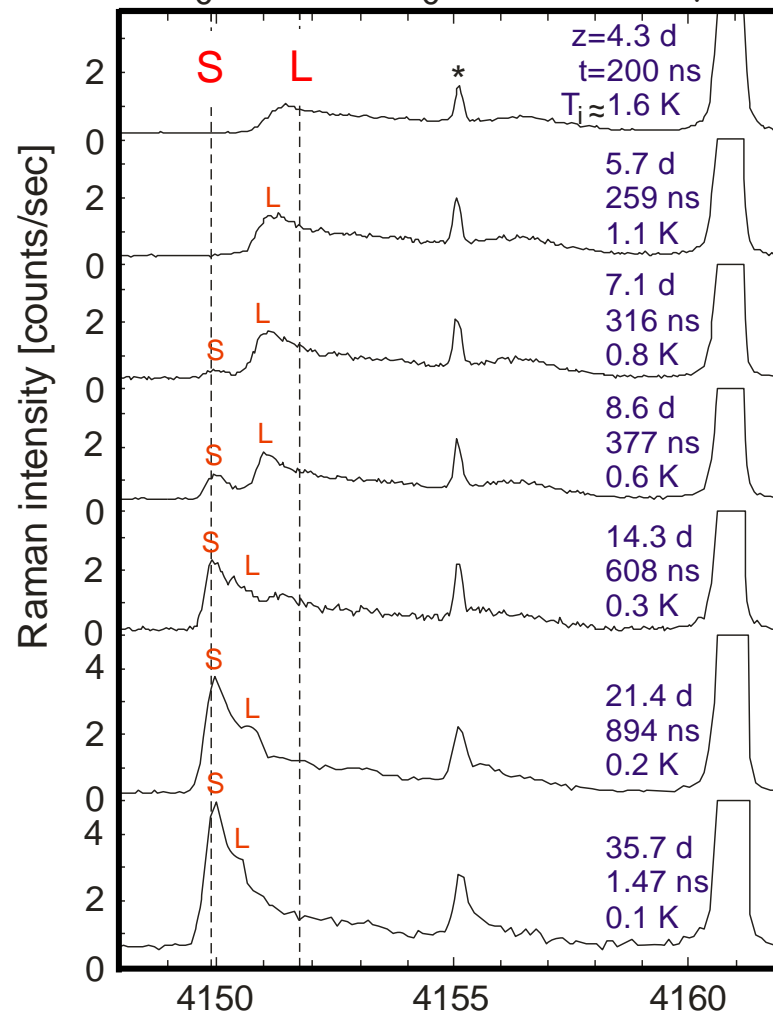
Magic Numbers Appear in High Resolution Raman Spectra of (pH₂) Clusters with Distance



T₀=46.5 K
P₀=1 bar
d=50 micron

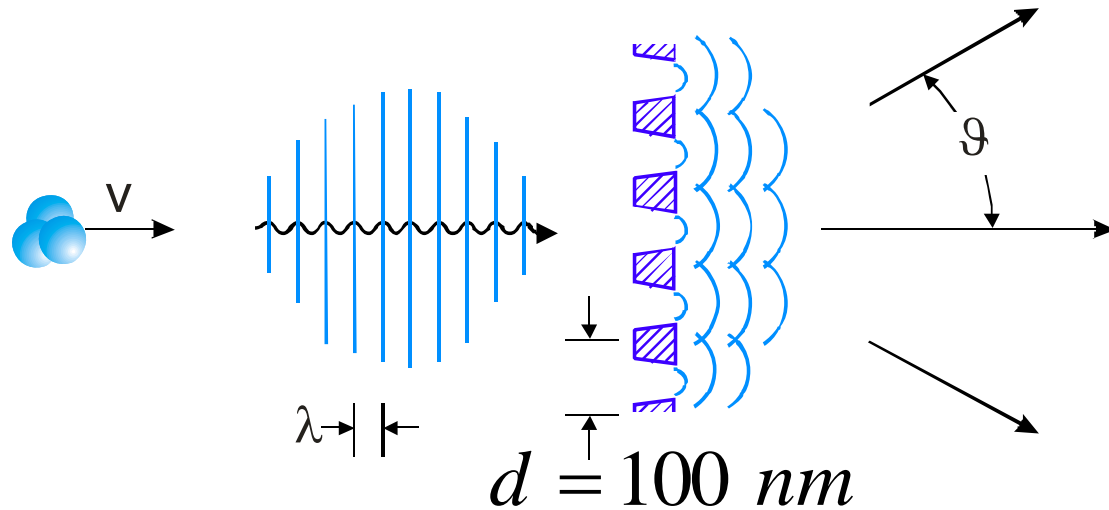
Evidence for a Liquid - Solid Transition in Large pH_2 - Clusters

$T_0 = 36.5 \text{ K}$, $P_0 = 2 \text{ bar}$, $d = 35 \mu$



Liquid Clusters also Recently Detected in He Seeded Beams:
Kuyanov-Przument and Vilesov, submitted

Non-destructive Transmission Grating Diffraction



N-particle Cluster
de Broglie Wavelength

$$\lambda = \frac{h}{N \times (mv)}$$

Typically: $\lambda = 0.5 \text{ \AA}$,

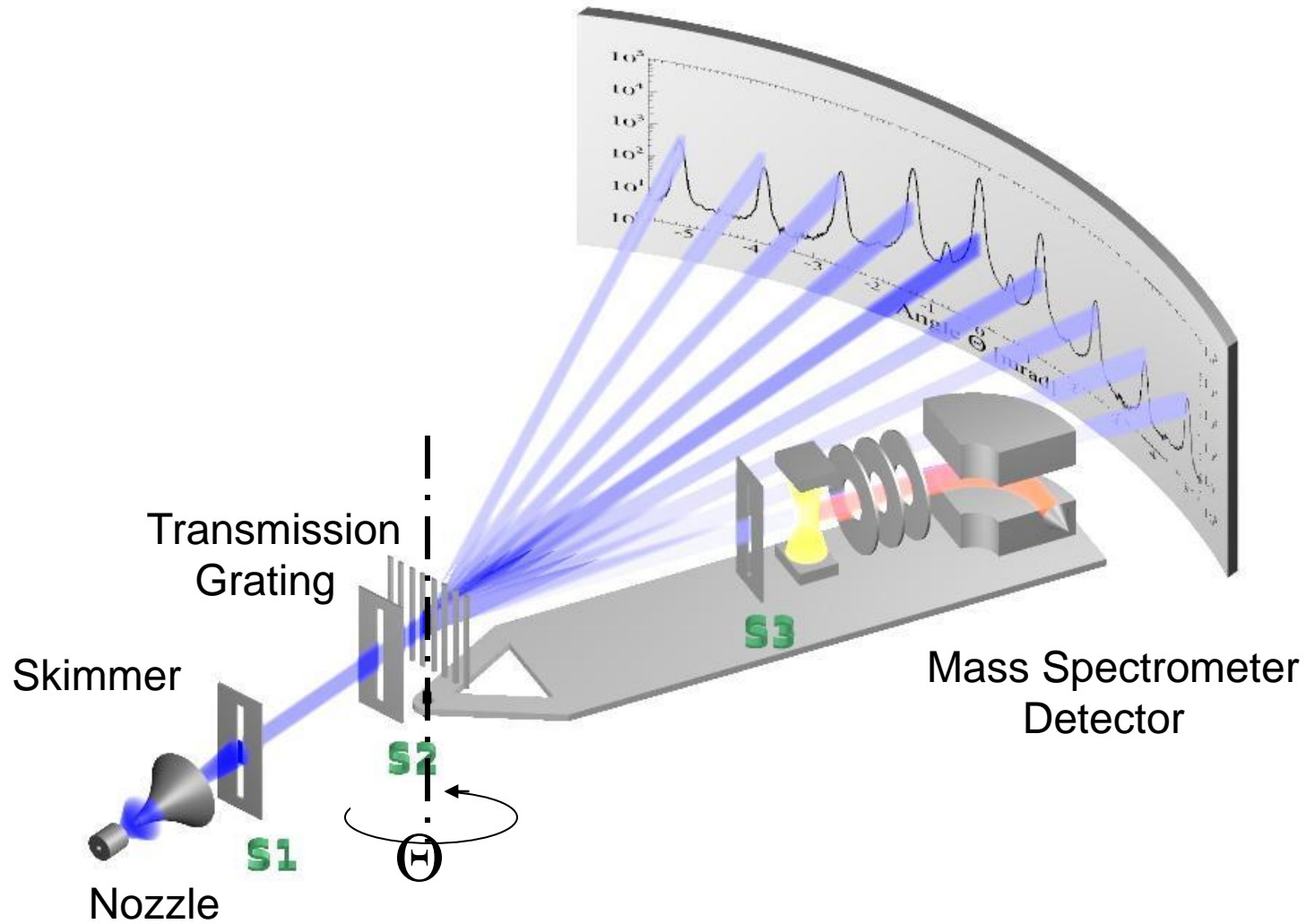
Diffraction Angles

$$\theta \approx \frac{\lambda}{d} \quad (n=1)$$

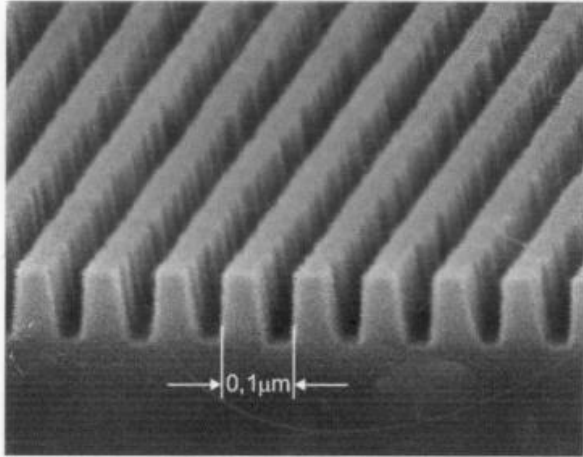
$\theta = 0.5 \text{ mrad}$

The diffraction angle is inversely prop. to N

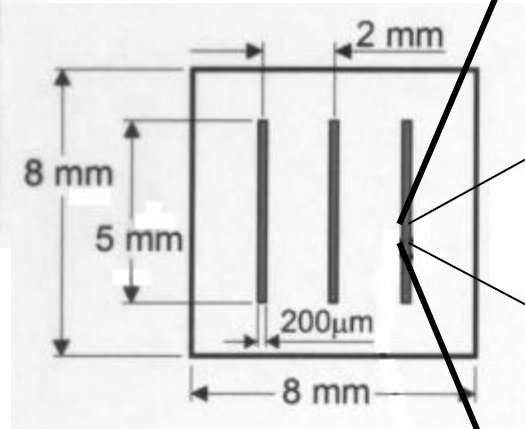
Schematic Diagram of the Diffraction Apparatus



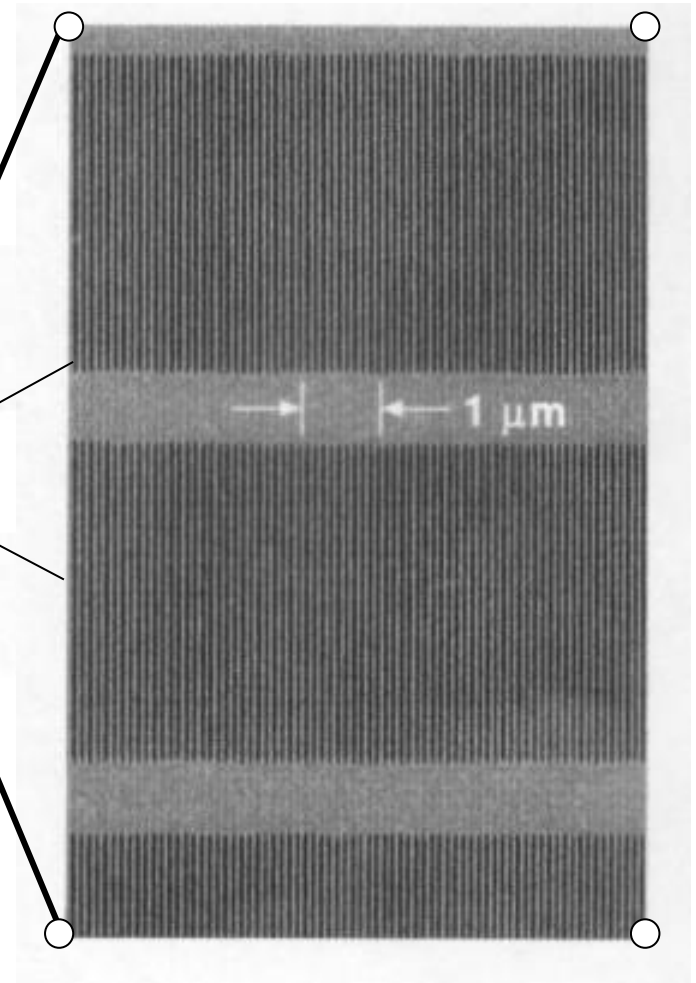
Electron Microscope Picture of the SiN_x Transmission Gratings



View of the chip before the Si base is etched away



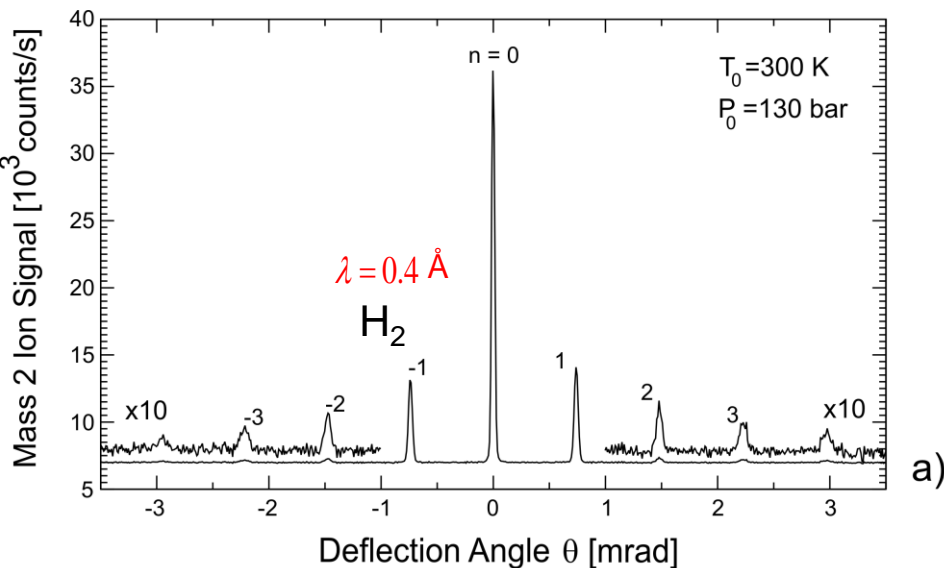
100 nm period
~50 nm slit width



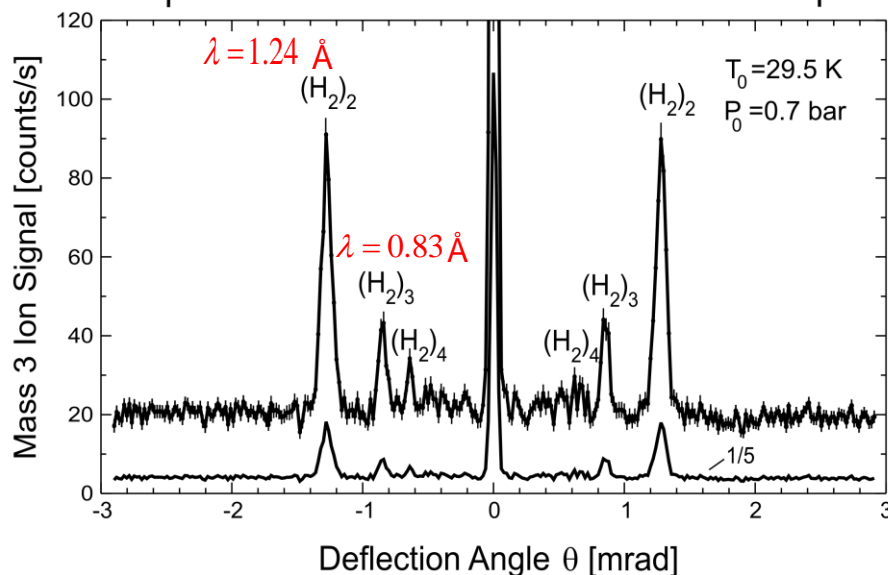
Courtesy of Prof. H. Smith and Dr. Tim Savas ,M. I. T.

Diffraction of Small Normal-H₂ Clusters

$P_0 = 130 \text{ bar}$,
 $T_0 = 300 \text{ K}$



$P_0 = 0.7 \text{ bar}$,
 $T_0 = 29.5 \text{ K}$



ionization

Kornilov and Toennies,
J. Chem.Phys. 128,194306 (2008)

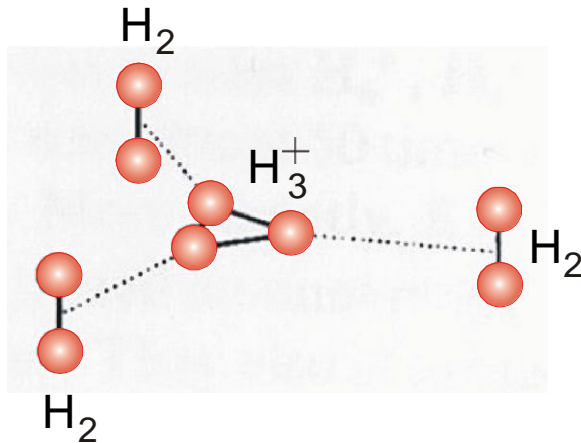
To Obtain the Neutral Cluster Size Distribution Need Fragmentation Pattern:

$$(H_2)_N + e^- \rightarrow H_K^+ + 2e^-$$

Only fragments with odd masses are present:

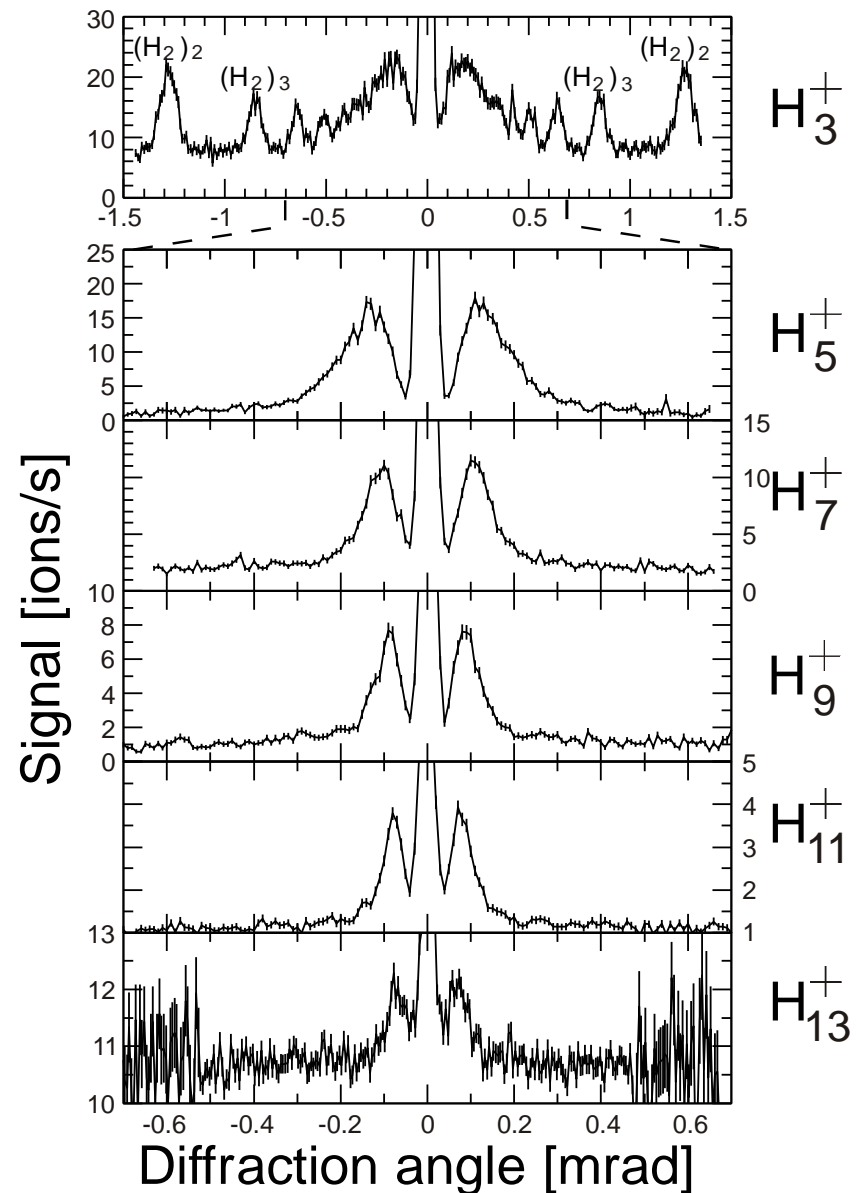
reaction $H_2^+ + H_2 \rightarrow H_3^+ + H$ has very large cross section.

Structure of ions:

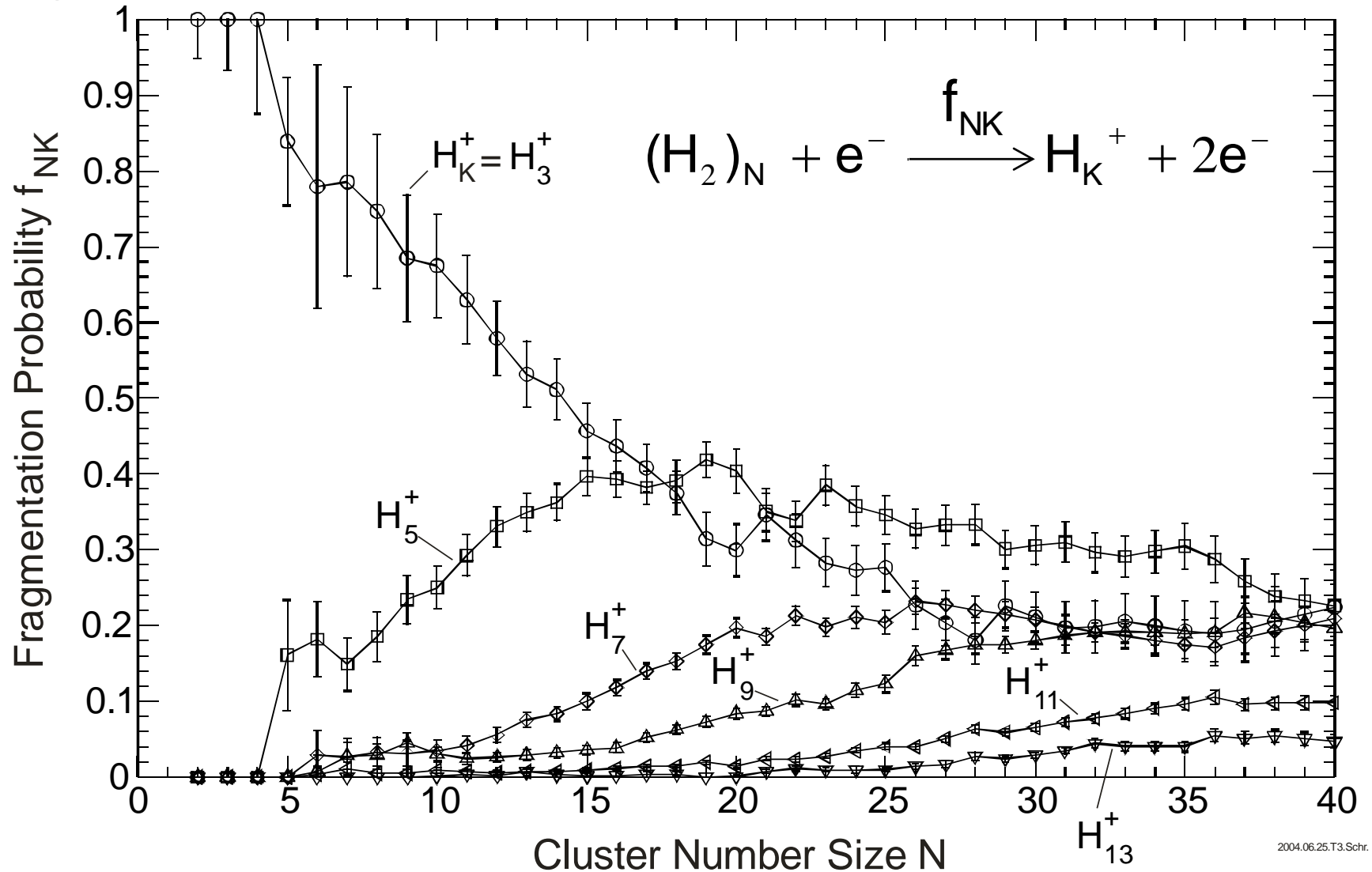


Kornilov and Toennies,
 J. Chem.Phys.128,194306 (2008)

f_{NK}

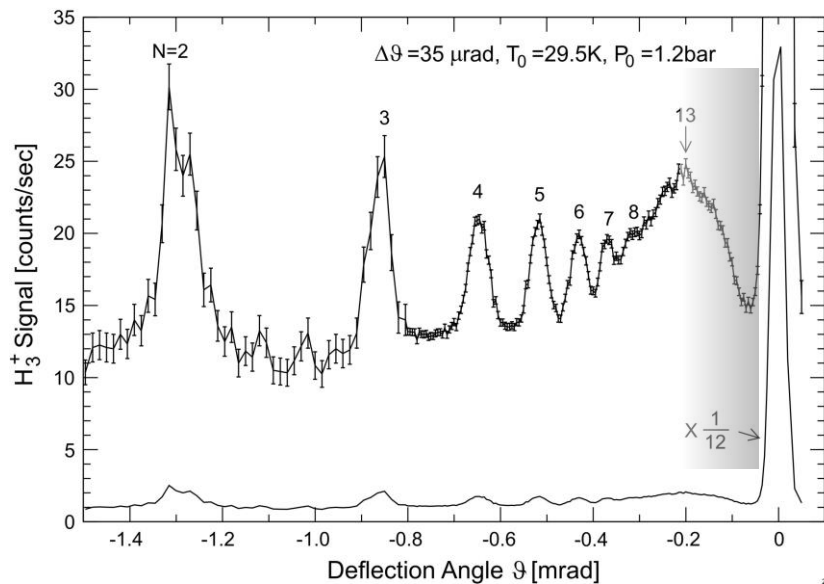


Fragmentation Probabilities as a Function of Cluster Size N



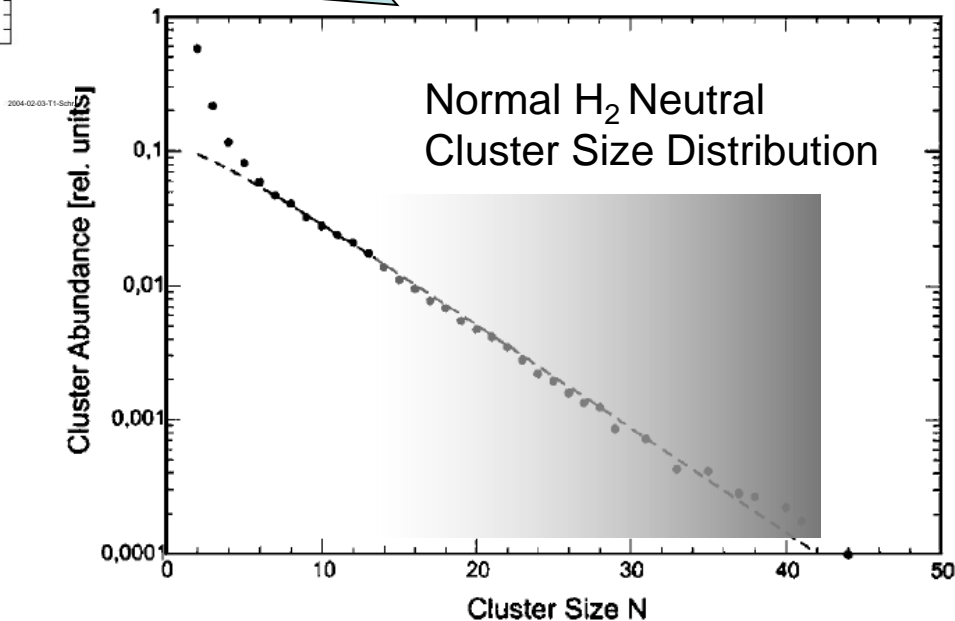
O. Kornilov and J. P. Toennies, in preparation

Determination of Neutral Cluster Size Distributions



$$G(N) = \frac{I_m(\theta)}{f_{NK} \cdot \sigma_{ion}} \left| \frac{\theta}{N^2} \right|$$

No Magic Numbers Seen
with normal H_2 !!



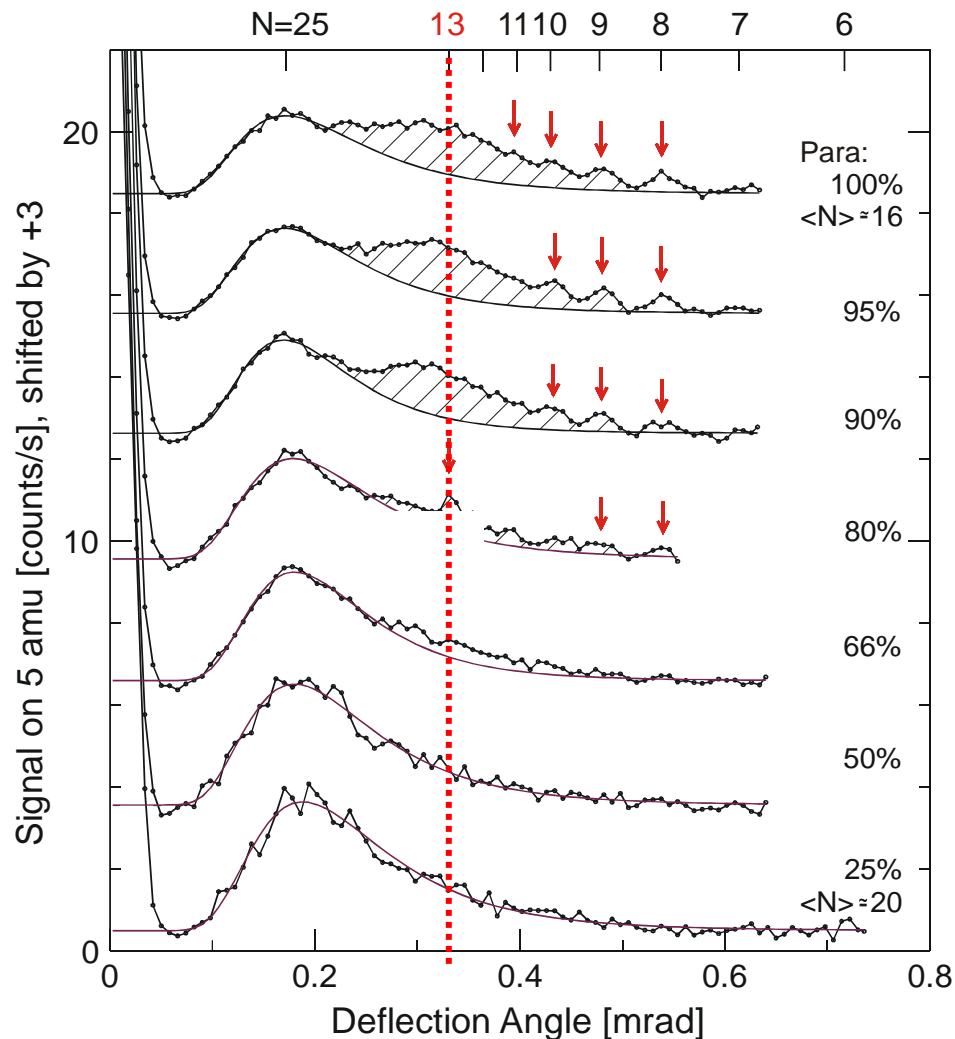
Magic no. with pH₂

Kornilov and Toennies, J. Chem.Phys. 128, 194306 (2008)

Only para-H₂ Exhibits Magic Number 13

5% H₂ in He,
T₀=20K, P₀=3.5 bar

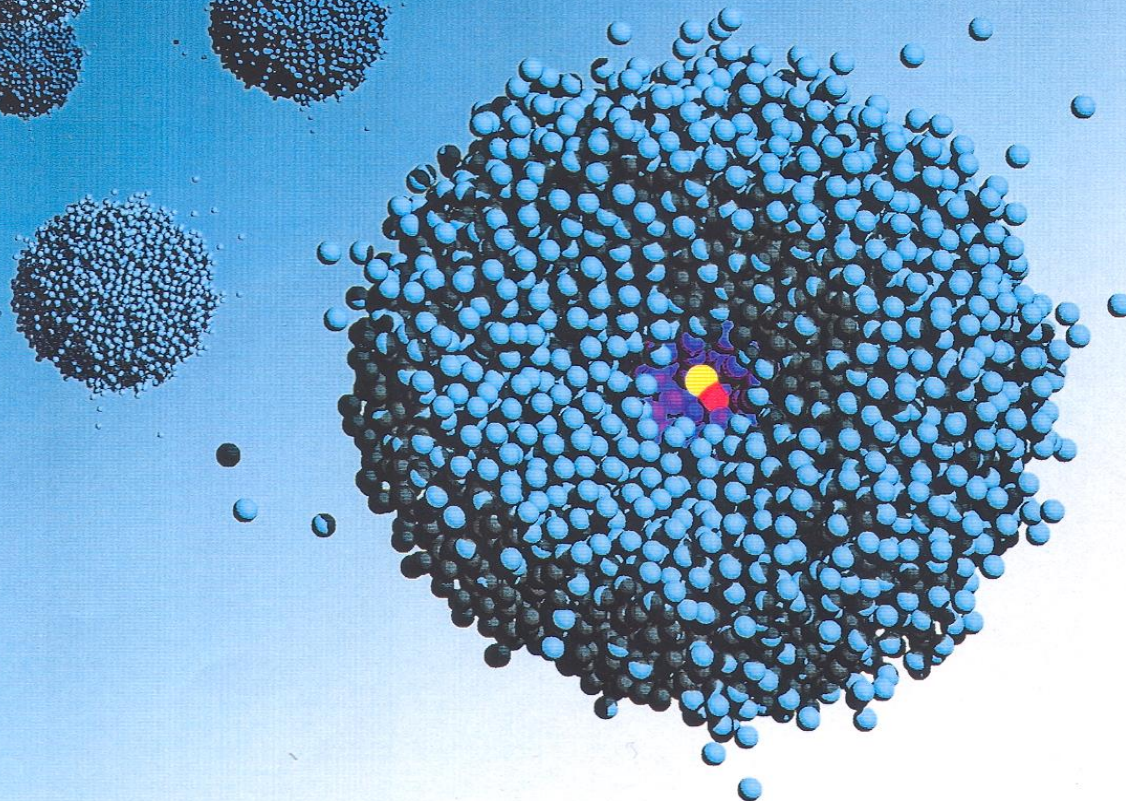
T_{int} ≤ 1.5 K



2004-09-03-T1-Schr.

Recent Calculations Indicate that Superfluidity is Suppressed by 3 or more Ortho-H₂ Molecules

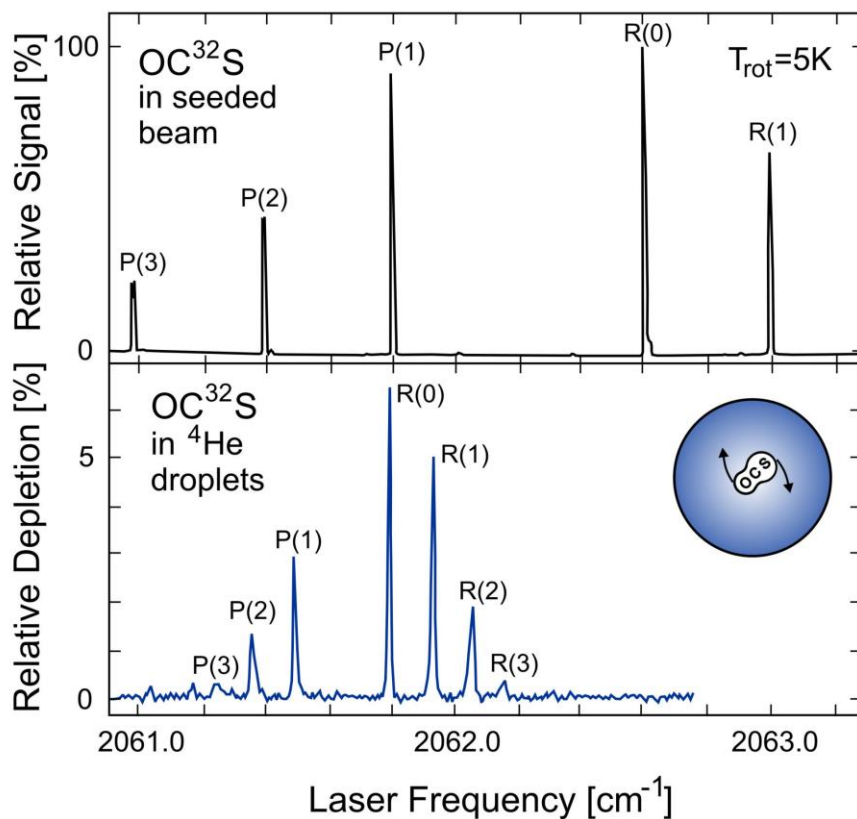
Droplet Spectroscopy



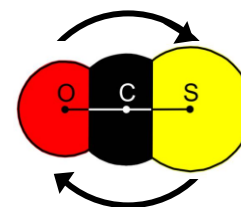
Rot in dropl

The Infra-red Spectra of Chromophore Molecules Inside Helium Droplets Indicate that the Molecules Rotate Freely.

A Microscopic Manifestation of Superfluidity



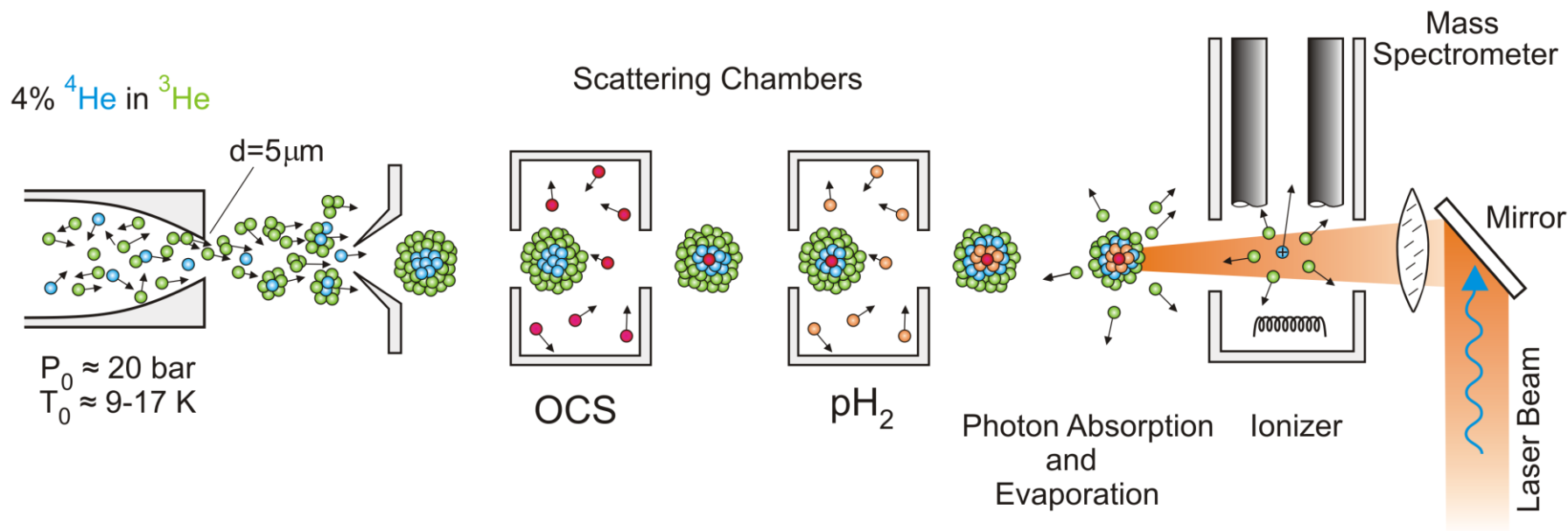
Boltzmann Plot of Intensities yields $T_{\text{Drop}}=0.37\text{K}$



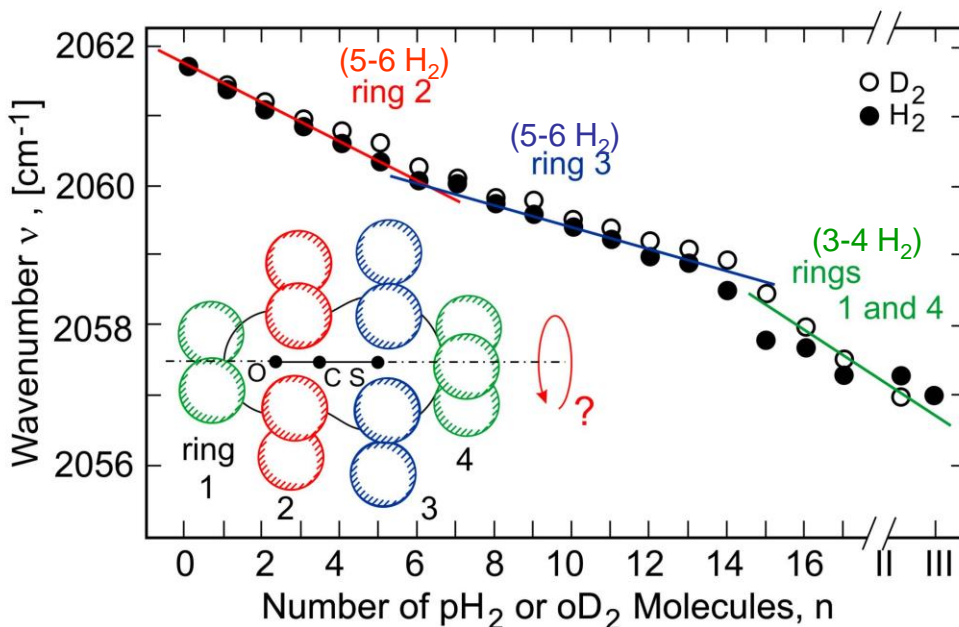
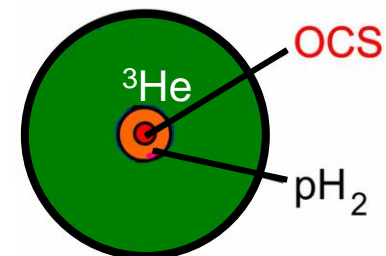
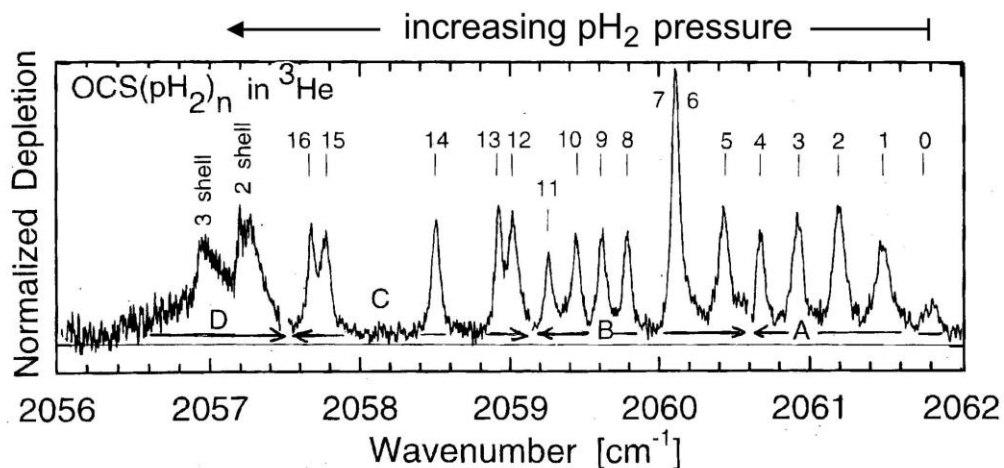
Sharp spectral features indicate that the molecule rotates without friction

The closer spacing of the lines indicates a factor 2.7 larger moment of inertia

OCS-(pH₂)_N Clusters are Assembled Inside ⁴He Droplets or Inside Mixed ⁴He/³He Droplets

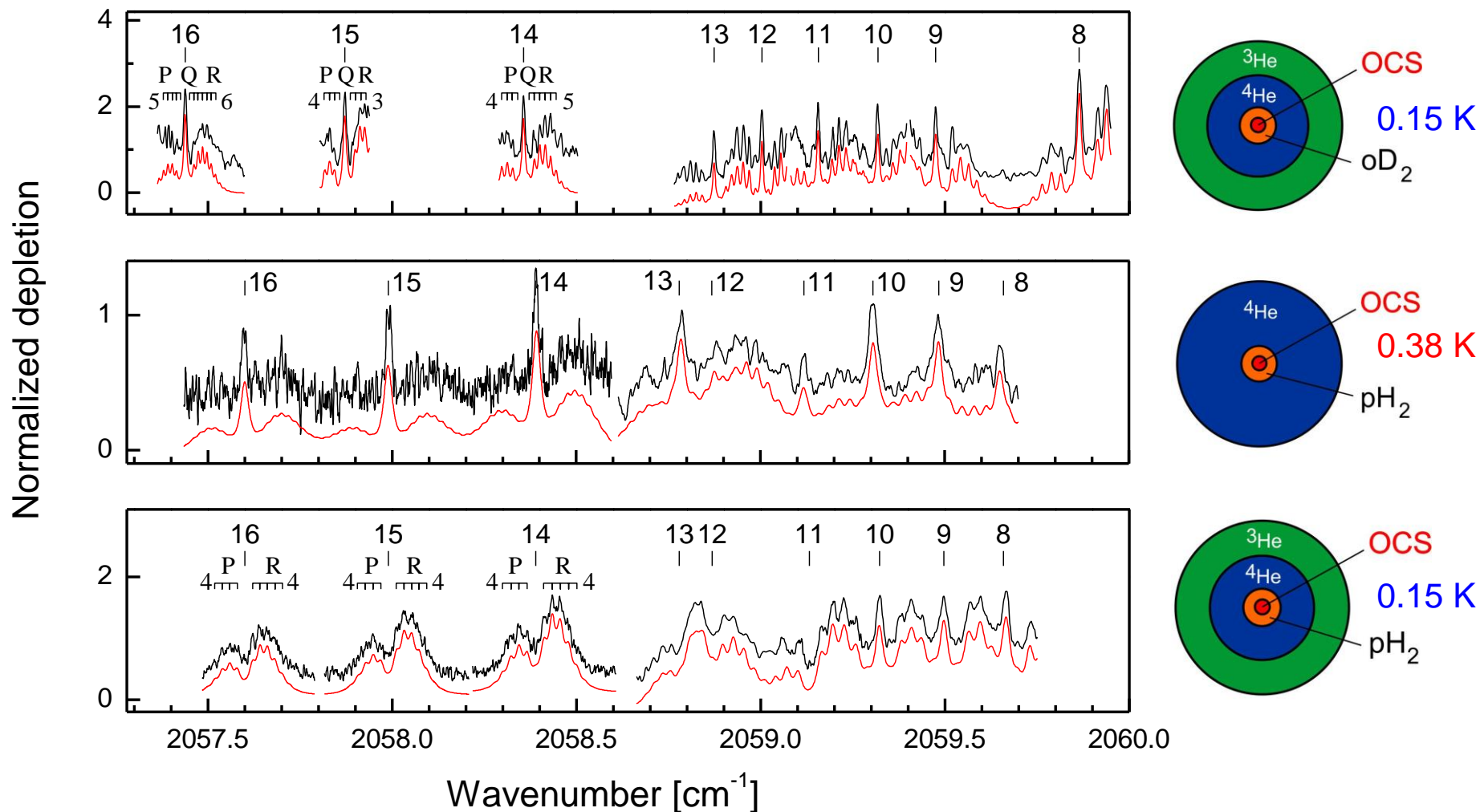


The Line Shifts of Resolved Cluster Peaks in Pure ^3He Droplets Indicate Ring Growth



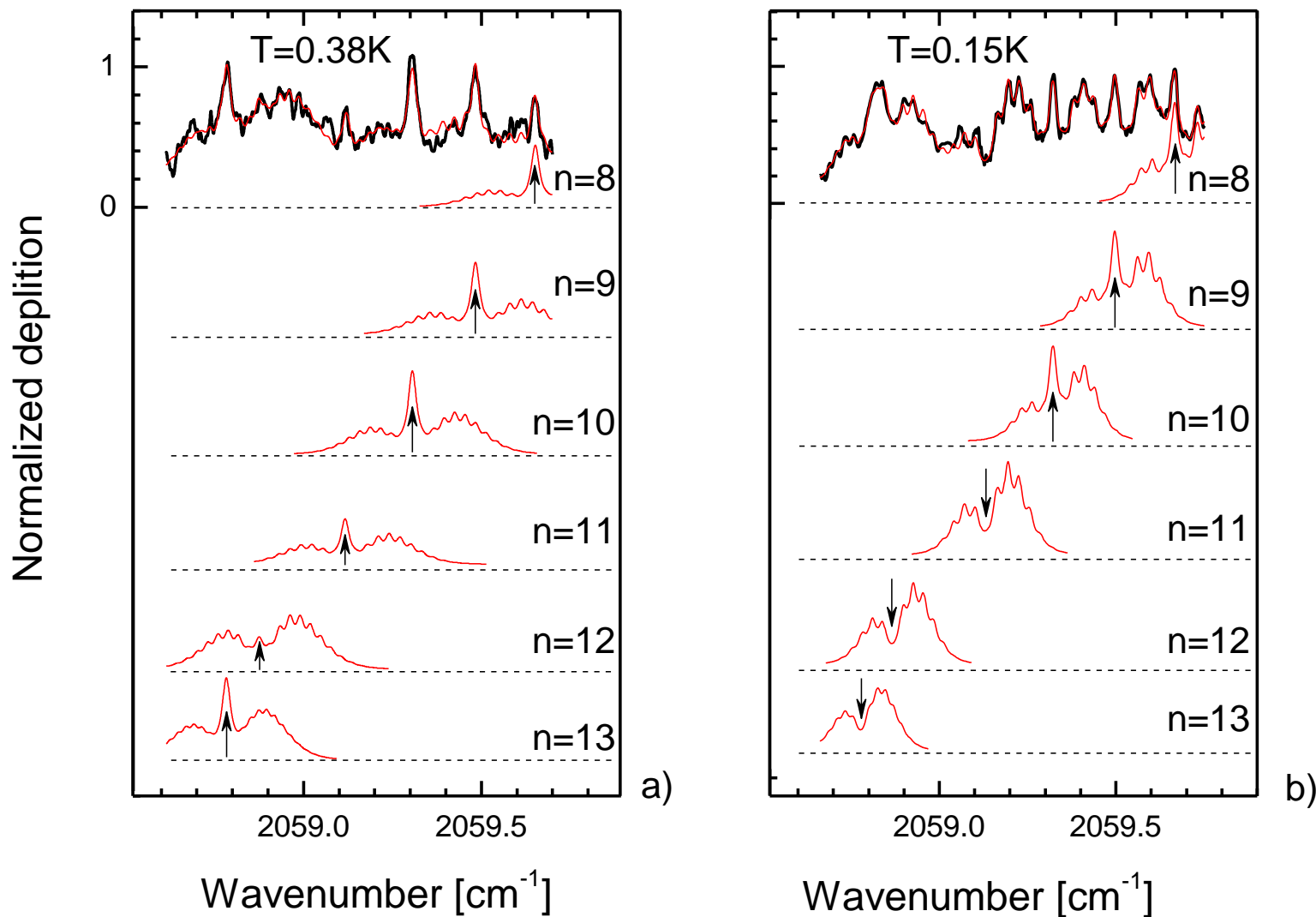
Q-branch

On Cooling from 0.38 to 0.15 K the Q-Branch remains for the o-D₂ Clusters but Vanishes for the p-H₂ Clusters

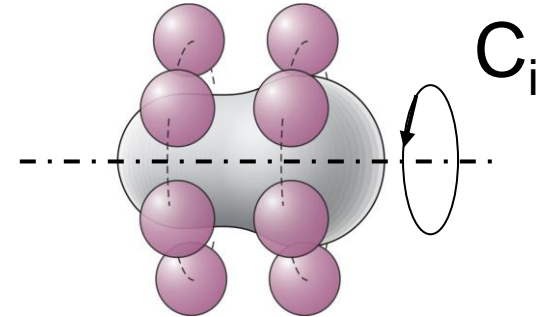
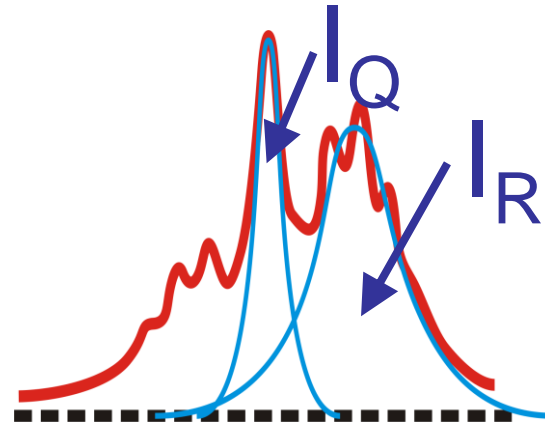


The Q-branch indicates that the cluster rotates about its symmetry axis

The Deconvolution of the Spectra Reveals that the Anomaly Starts at $n \geq 11$ (except $n = 12$)



The Spectrum is Characterized by the relative Areas of the Q- and R-Branched



$$I_Q \propto \sum_{JK} F(K, J) g_{JK} w_{JK} \exp[-E_{rot}(J, K) / kT]$$

$$I_R \propto \sum_{JK} G(K, J) g_{JK} w_{JK} \exp[-E_{rot}(J, K) / kT]$$

where for distinguishable particles $w_{JK} = 1$

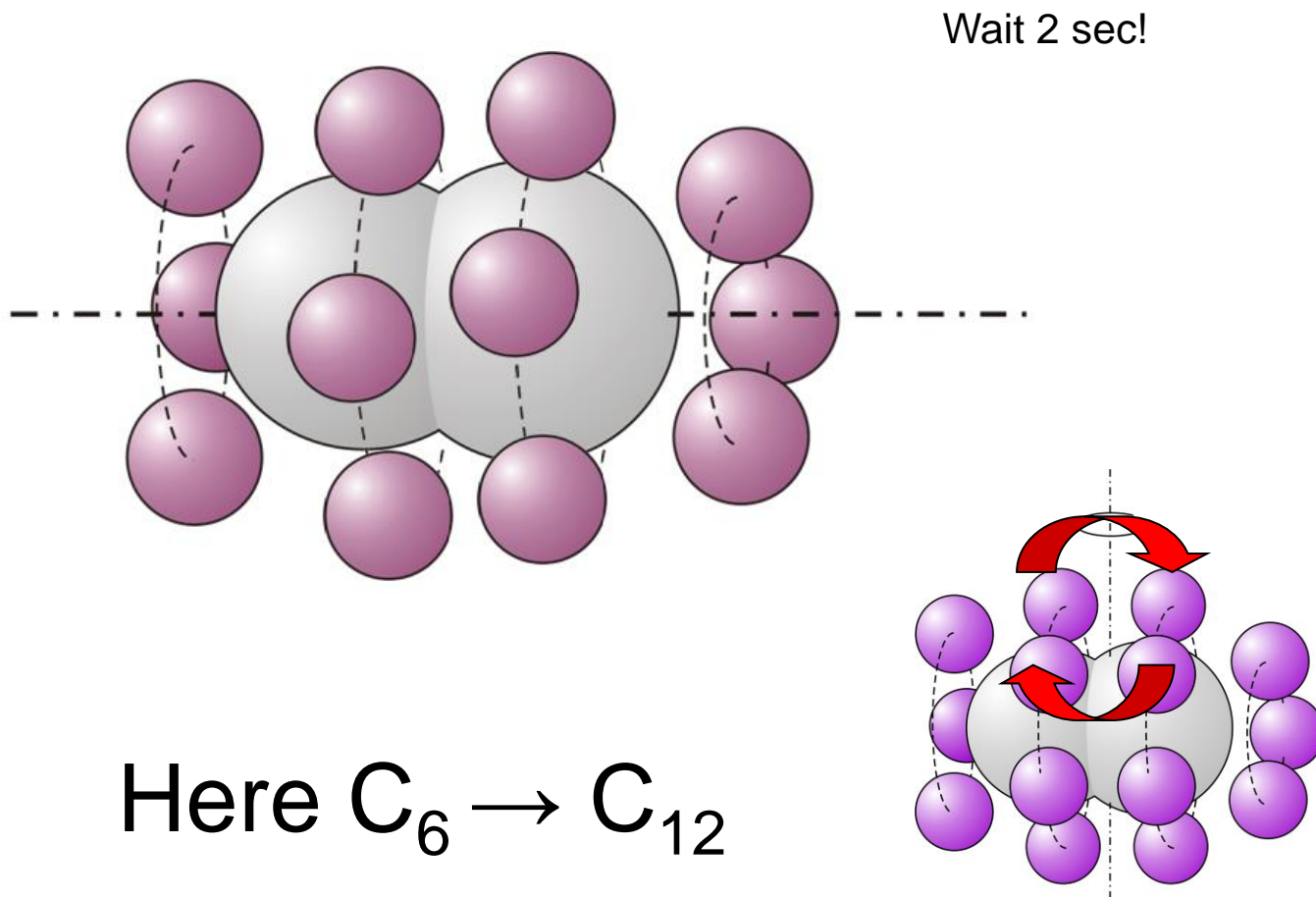
indistinguishable particles $w_{JK} = \delta_{k,m \cdot i} \quad m = 0, 1, 2, \dots$

The experimental ratio I_Q/I_R depends on the axial symmetry of the cluster C_i

Drop in I_Q/I_R at $T = 0.15$ K for $N \geq 11$ (except. 12)
 Corresponds to Large Increase in Symmetry i

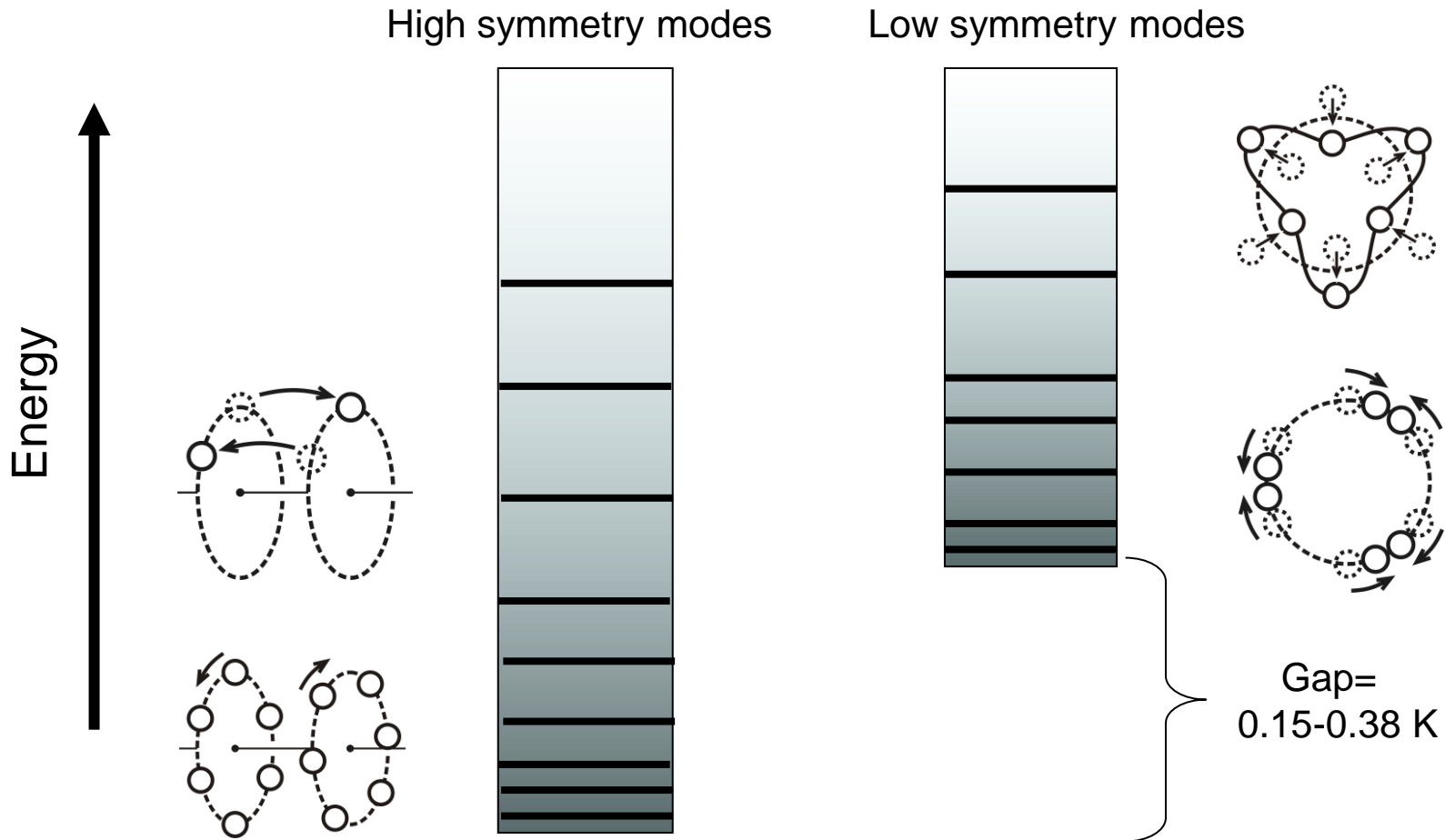
N	T = 0.38 K	i	T = 0.15 K	i	
10	0.56 ± 0.05	≤ 2	0.48 ± 0.04	1	same
11	0.51 ± 0.17	≤ 3	≤ 0.09	≥ 5	increase
12	≤ 0.19	≥ 6	≤ 0.07	≥ 4	same
13	0.62 ± 0.24	≤ 3	≤ 0.07	≥ 4	increase
14	0.69 ± 0.15	≤ 2	≤ 0.07	≥ 4	increase
15	0.64 ± 0.20	≤ 3	≤ 0.07	≥ 5	increase
16	0.62 ± 0.20	≤ 3	≤ 0.07	≥ 5	increase

From the Ethane Molecule it is known that
Hindered Rotations Double the Axial Symmetry



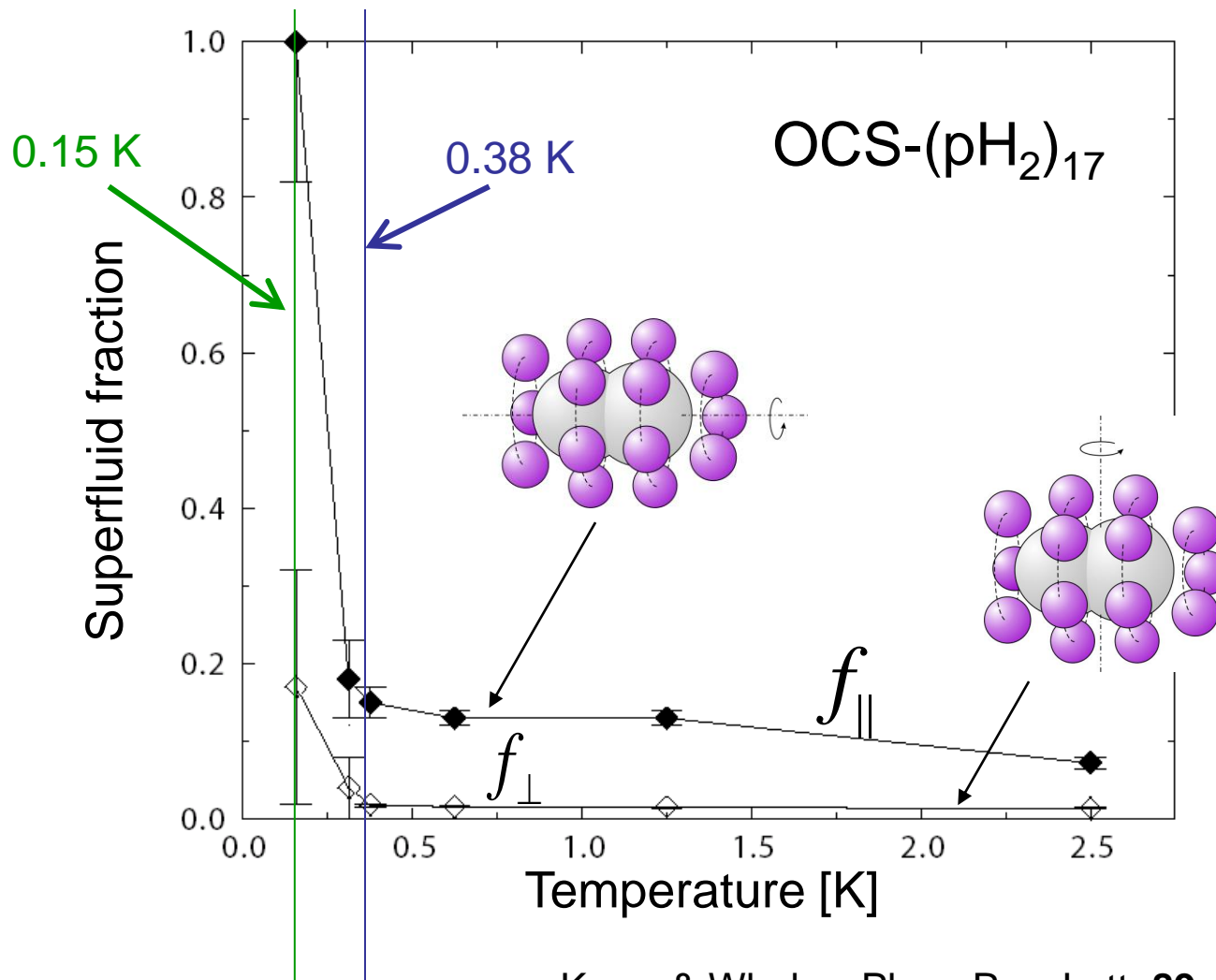
Energy levels

Postulated Energy Level Diagram



Is this Really Superfluidity??

Yes, We Think So Since the PIMC Calculations also Predict a Superfluid Axial Response at About the Same Temperature



concl

Summary

Small H₂ Clusters (N≤26) are fascinating objects exhibiting superfluidity while at the same appearing to be solid

The few experiments confirm the existence of magic numbers in pure para H₂ clusters, but not in normal H₂ clusters. Do magic nos. correlate with superfluidity??

Only evidence for superfluidity comes from experiments with doped clusters. There the effect is attributed to low frequency modes of high axial symmetry which facilitate particle exchanges

HYDROGEN I

This molecule hydrogen is so very small
Yet it is the most abundant of all

It only has atoms two,
And as many electrons too.

It's every theoreticians delight,
they think they can calculate it just right.

President George Bush says its great stuff
to fight global warming it is enough!

Even though Ginzburg had a theory convincing.
It's superfluidity is no simple thing!

HYDROGEN II

It seems not so difficult after all,
if clusters like the molecule are small

Doping with OCS provides a clue
But still there is much for us to do!

Thank You!

Please Ask Questions!

End

Please Ask Questions!